

# Protostellar collapse: protostar and disk formation

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# Outline

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## 1. Introduction

- Context
- What has changed since 2014?

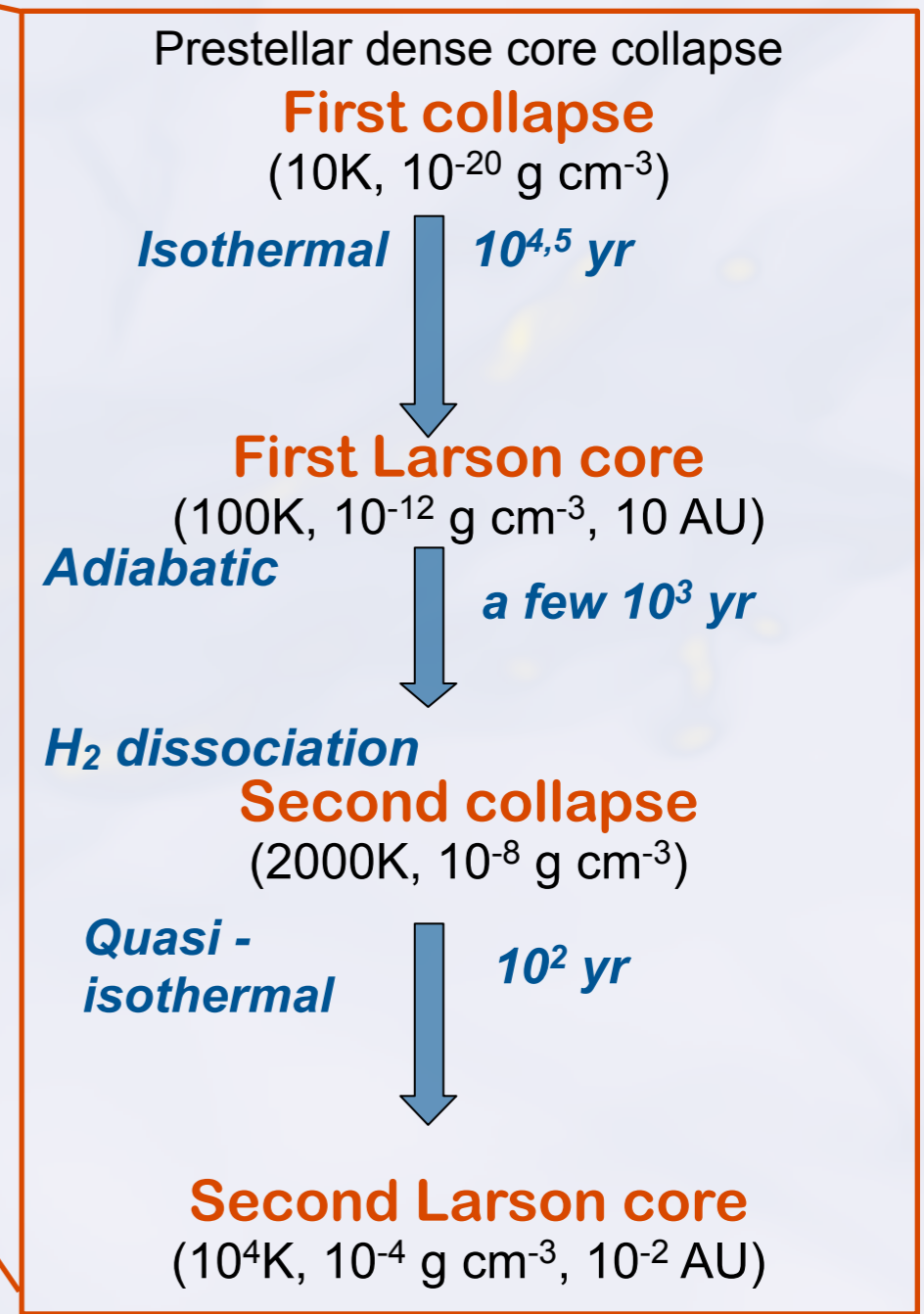
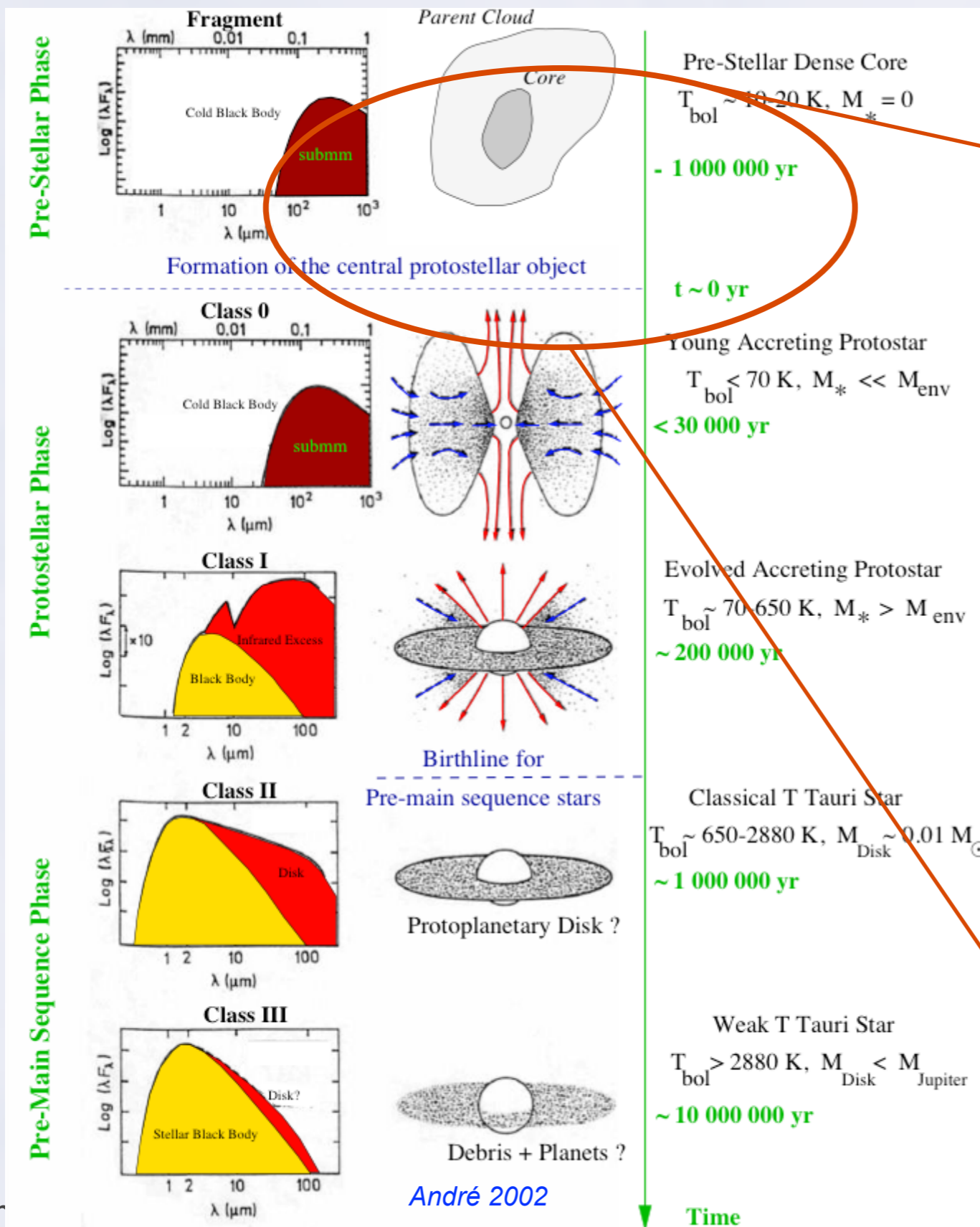
## 2. From dense core collapse....

- Disc formation
- Comparison with ALMA observations

## 3. To protostar and protoplanetary disc birth

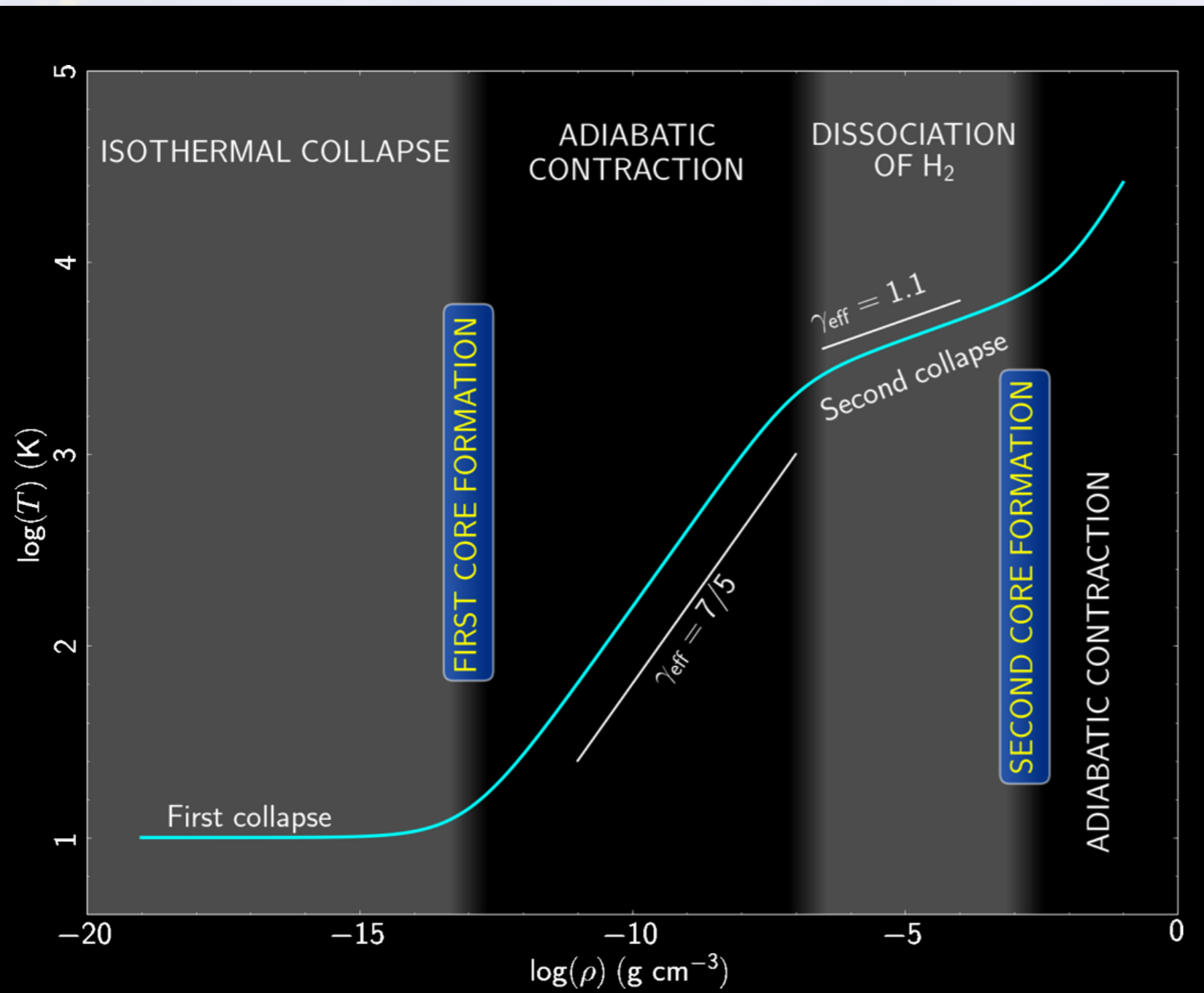
- Early evolution
- Dust and gas mixture

# Star formation evolutionary sequence

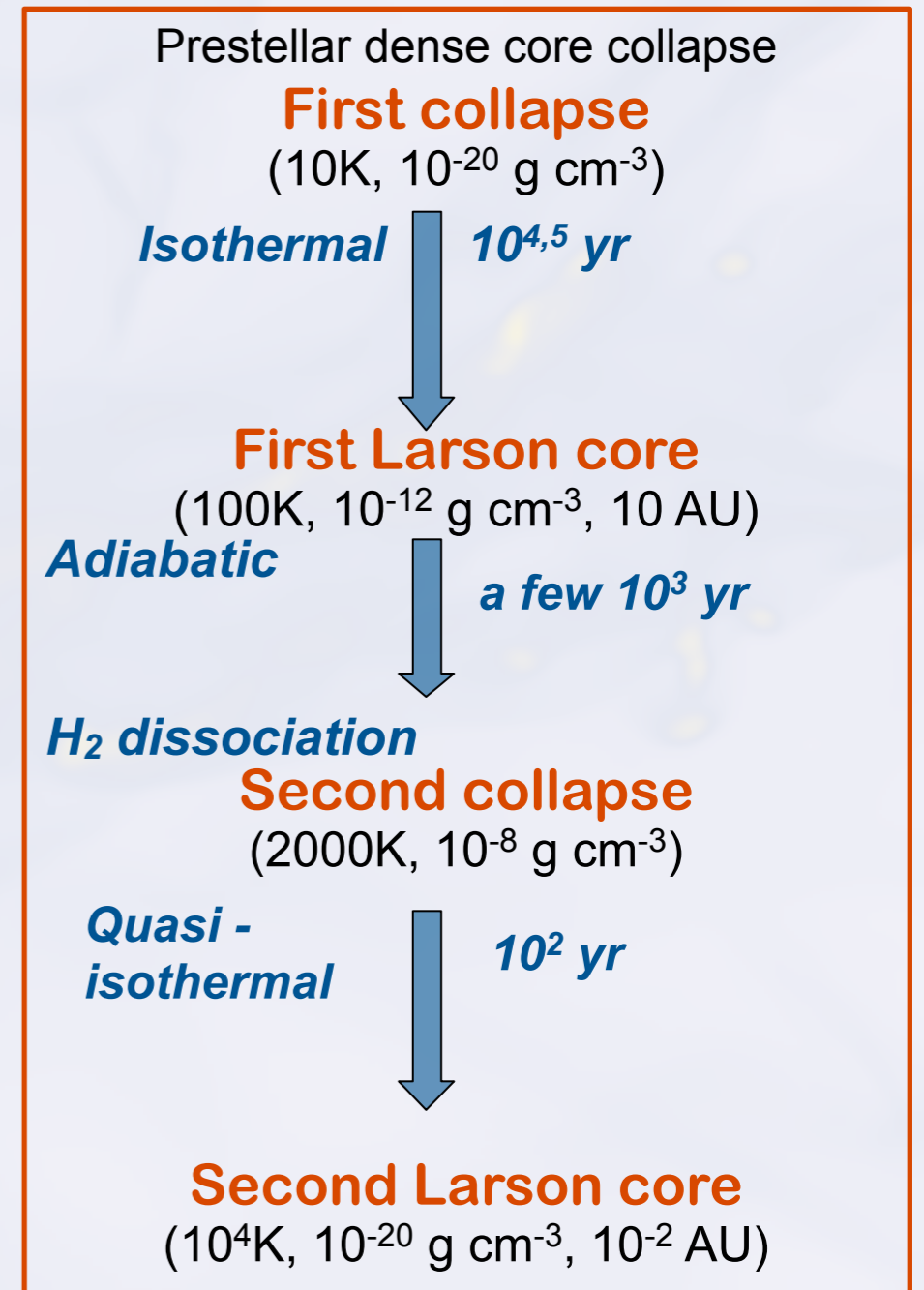




# Star formation evolutionary sequence



Vaytet et al. (2013)



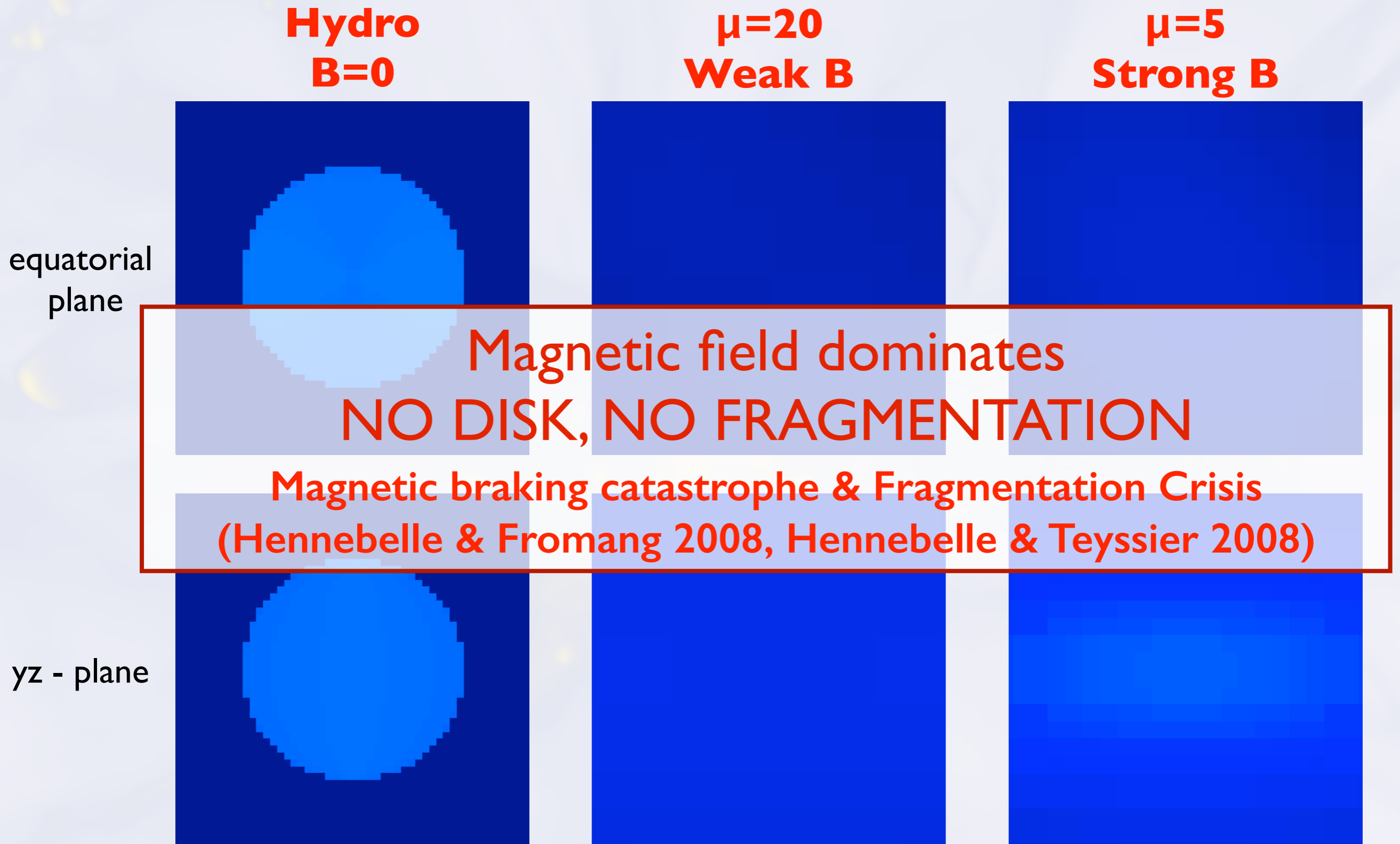


# Main achievements since 2014

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- **Observations by ALMA & NOEMA**
  - down to a spatial resolution  $< 50$  AU in Class 0 protostars  
=> the fragmentation and disk scales!
  - dust polarization  
=> magnetic fields
- **New numerical tools:**
  - non-ideal MHD : ambipolar diffusion, Ohmic diffusion, and Hall effect
  - multifrequency radiation hydrodynamics
  - dust and gas mixture

# State-of-the-art in 2010: ideal MHD



# Disk formation in magnetised cores

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## ✓ Late formation

- end of class 0,  $M_{\text{env}} \ll M_{\text{env},0}$  (e.g., [Machida & Hosokawa 2013](#))

## ✓ Misalignment

- no reason for the rotation axis and the magnetic field to be aligned (e.g., [Hull et al. 2013](#))
- reduces magnetic braking efficiency (e.g., [Hennebelle & Ciardi 2009](#), [Joos et al. 2012](#), [Li et al. 2013](#))

## ✓ Turbulent diffusion

- reconnection events fast with Ohmic diffusion only, collective effect at larger scale (e.g., [Santos Lima et al. 2012](#), [Joos et al. 2013](#), [Seifried et al. 2013](#))

## ✓ Non-ideal MHD

- Ohm dissipation ([Tomida et al. 2013, 2015](#), [Machida et al.](#))
- Hall effect ([Krasnopolsky et al. 2011](#), [Tsukamota et al. 2015](#), [Wurster et al. 2016](#))
- ambipolar diffusion ([Tsukamota et al. 2015](#), [Masson et al. 2016](#))



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- ambipolar diffusion (*Tsukamota et al. 2015, Masson et al. 2016*)

# Non-ideal MHD

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$$\frac{\partial \mathbf{B}}{\partial t} - \nabla \times \left[ \mathbf{u} \times \mathbf{B} - \eta_{\Omega} \mathbf{J} - \frac{\eta_H}{\|\mathbf{B}\|} \mathbf{J} \times \mathbf{B} + \frac{\eta_{AD}}{\|\mathbf{B}\|^2} \mathbf{J} \times \mathbf{B} \times \mathbf{B} \right] = 0$$

## Non-ideal effects:

- rearrangement of magnetic field lines
- reconnection
- magnetic flux diffusion
- ... needs gas-grain chemistry

## ✓ Non-ideal MHD

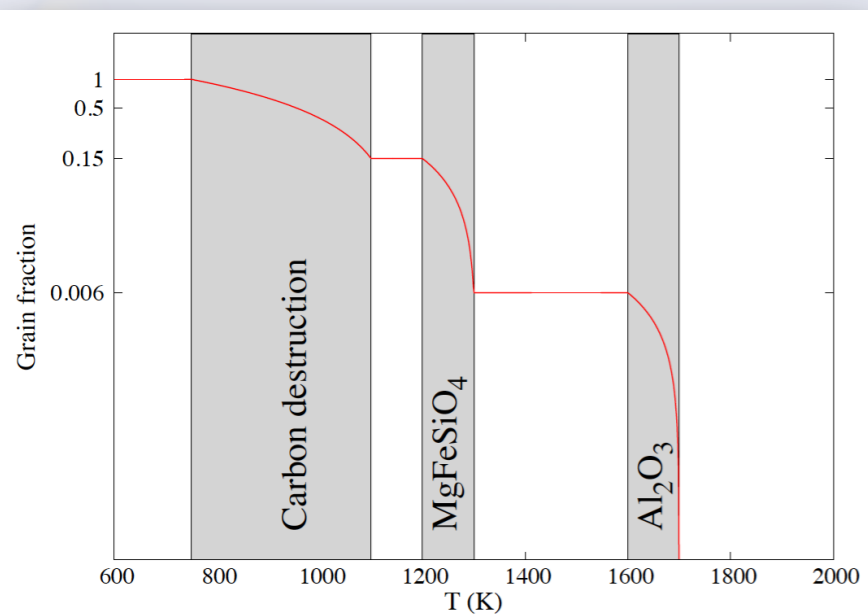
- Ohm dissipation (*Tomida et al. 2013, 2015, Machida et al.*)
- Hall effect (*Krasnopolsky et al. 2011, Tsukamota et al. 2015, Wurster et al. 2016*)
- ambipolar diffusion (*Tsukamota et al. 2015, Masson et al. 2016*)

# Equilibrium chemistry for non-ideal MHD

## ✓ Reduced chemical network dedicated for ionisation (based on the work by *Umebayashi & Nakano 1990*)

- H, He, C, O, metallic elements (Fe, Na, Mg, etc..)
- $H^+$ ,  $H_3^+$ ,  $He^+$ ,  $C^+$ , molecular and metallic ions
- bins in the dust grains size distribution ( $G$ ,  $G^+$ ,  $G^-$ )
- dust evaporation at  $T > 800$  K
- thermal ionisation of potassium ( $T > 1000$  K)
- neutral elements have constant abundances

Reaction	$\alpha$	$\beta$	$\gamma$
$H^+ + O \rightarrow H + O^+$	$6.86 \times 10^{-10}$	0.26	0
$H^+ + O_2 \rightarrow H + O_2^+$	$2.00 \times 10^{-9}$	0.00	0
$H^+ + M \rightarrow H + M^+$	$1.10 \times 10^{-9}$	0.00	0
$He^+ + H_2 \rightarrow He + H^+ + H$	$3.70 \times 10^{-14}$	0.00	35
$He^+ + CO \rightarrow He + C^+ + O$	$1.60 \times 10^{-9}$	0.00	0
$He^+ + O_2 \rightarrow He + O^+ + O$	$1.10 \times 10^{-9}$	0.00	0
$H_3^+ + CO \rightarrow H_2 + HCO^+$	$1.36 \times 10^{-9}$	-0.14	0
$H_3^+ + O \rightarrow H_2 + OH^+$	$7.98 \times 10^{-10}$	-0.16	0
$H_3^+ + O_2 \rightarrow H_2 + O_2H^+$	$9.30 \times 10^{-10}$	0.00	0
$H_3^+ + M \rightarrow H_2 + H + M^+$	$1.10 \times 10^{-9}$	0.00	0
$C^+ + H_2 \rightarrow CH_2^+ + hv$	$2.00 \times 10^{-16}$	0.00	0
$C^+ + O_2 \rightarrow CO^+ + O$	$3.42 \times 10^{-10}$	0.00	0
$C^+ + O_2 \rightarrow CO + O^+$	$4.54 \times 10^{-10}$	0.00	0
$C^+ + M \rightarrow C + M^+$	$1.10 \times 10^{-9}$	0.00	0
$m^+ + M \rightarrow m + M^+$	$2.90 \times 10^{-9}$	0.00	0
$H^+ + e^- \rightarrow H + hv$	$3.50 \times 10^{-12}$	-0.75	0
$He^+ + e^- \rightarrow He + hv$	$5.36 \times 10^{-12}$	-0.5	0
$H_3^+ + e^- \rightarrow H + H + H$	$2.34 \times 10^{-8}$	-0.52	0
$H_3^+ + e^- \rightarrow H_2 + H$	$2.34 \times 10^{-8}$	-0.52	0
$C^+ + e^- \rightarrow C + hv$	$2.36 \times 10^{-12}$	-0.29	0
$m^+ + e^- \rightarrow m_1 + m_2$	$2.40 \times 10^{-7}$	-0.69	0
$M^+ + e^- \rightarrow M + hv$	$2.78 \times 10^{-12}$	-0.68	0
$H_2 \rightarrow H_2^+ + e^-$	$1.2 \times 10^{-17}$		
$H_2 \rightarrow H^+ + H + e^-$	$2.86 \times 10^{-19}$		
$He \rightarrow He^+ + e^-$	$6.58 \times 10^{-18}$		



- ✓UMIST database for gas species (*McElroy et al. 2013*)
- ✓Kunz & Mouschovias (2009) for interactions with and between grains

## ✓ Goal: compute a 3D table of abundances:

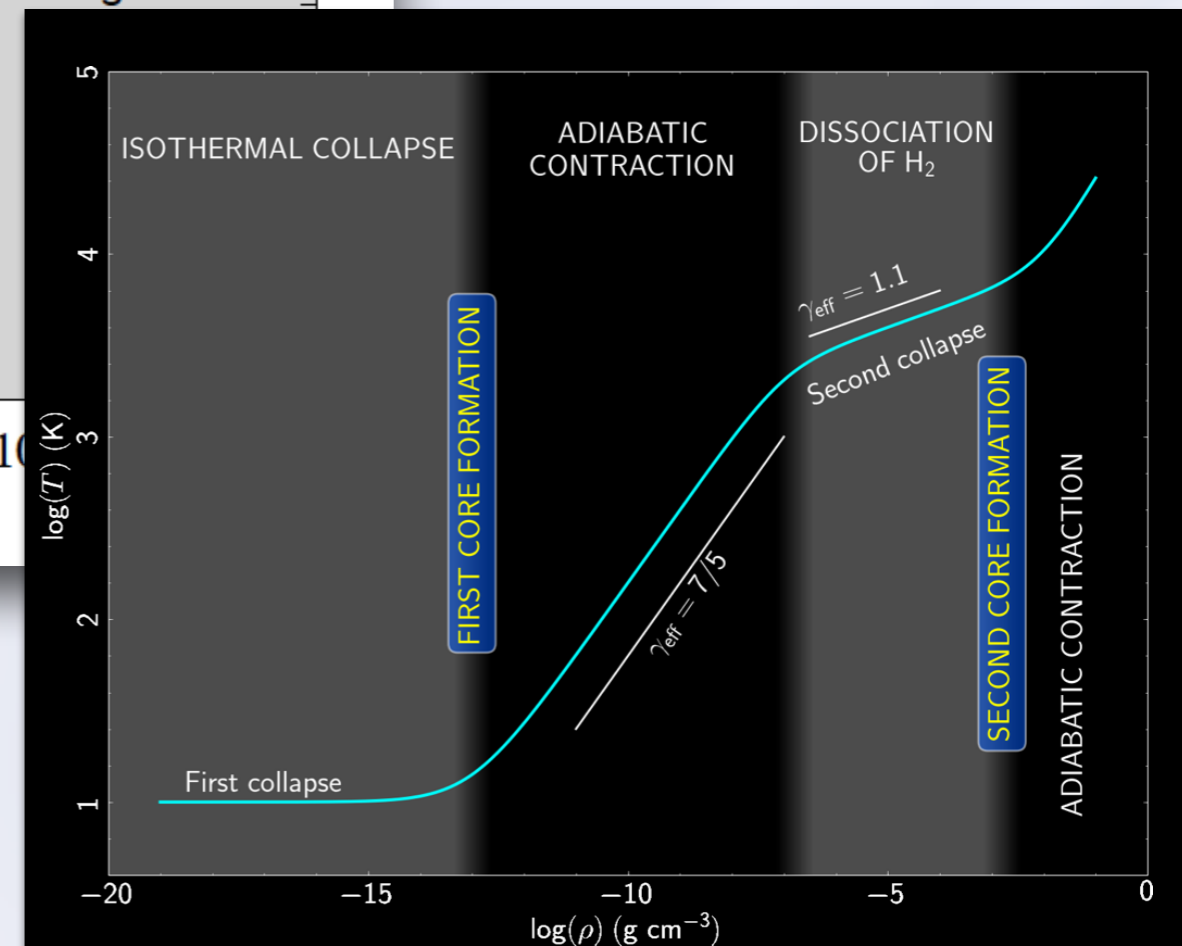
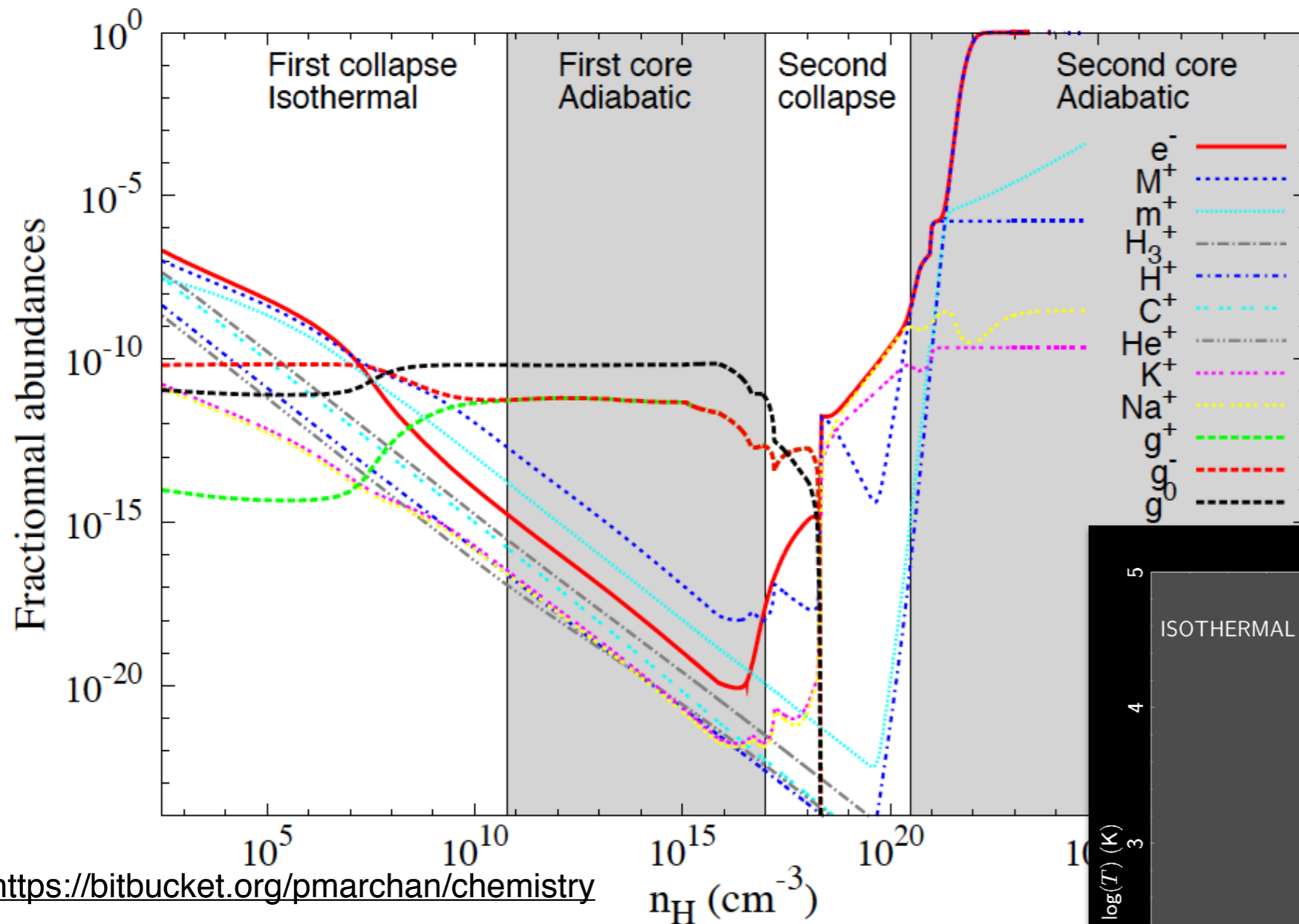
- depends on temperature, density and CR ionisation
- used on-the-fly in 3D calculations to compute resistivities

*Marchand et al. (2016)*

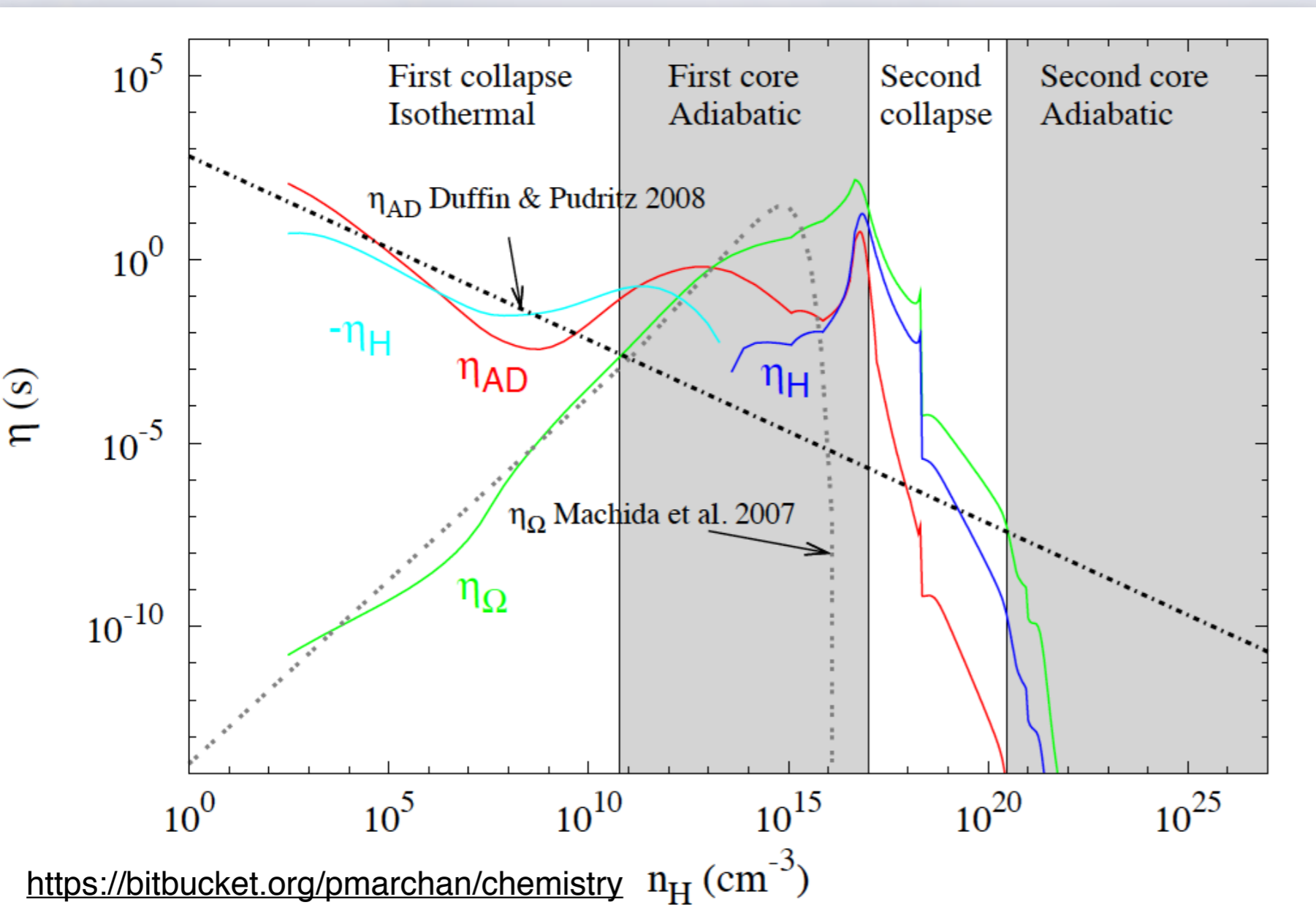


# Equilibrium chemistry for non-ideal MHD: results

*Marchand et al. (2016)*



# Equilibrium chemistry for non-ideal MHD: results



$$\eta_{\Omega} = \frac{1}{\sigma_{\parallel}},$$

$$\eta_{\text{H}} = \frac{\sigma_{\text{H}}}{\sigma_{\perp}^2 + \sigma_{\text{H}}^2},$$

$$\eta_{\text{AD}} = \frac{\sigma_{\perp}}{\sigma_{\perp}^2 + \sigma_{\text{H}}^2} - \frac{1}{\sigma_{\parallel}},$$

$$\sigma_{\parallel} = \sum_i \sigma_i,$$

$$\sigma_{\perp} = \sum_i \frac{\sigma_i}{1 + (\omega_i \tau_{\text{in}})^2},$$

$$\sigma_{\text{H}} = - \sum_i \frac{\sigma_i \omega_i \tau_{\text{in}}}{1 + (\omega_i \tau_{\text{in}})^2}.$$

- 1/ **Grain evaporation** is the most important effect
- 2/ Needs at least 10 bins in dust grain size distribution to converge...

*Marchand et al. (2016)*

# Radiation-magneto-hydrodynamics in **RAMSES**

- ✓ **Adaptive-mesh-refinement** code **RAMSES** (Teyssier 2002)
- ✓ **Non-ideal MHD** solver using Constrained Transport (Teyssier et al. 2006, Fromang et al. 2006, Masson et al. 2012, 2016). Resistivity from **steady-state gas-grain** chemistry (Marchand et al. 2016)
- ✓ **Multifrequency Radiation-HD** solver using the **Flux Limited Diffusion** approximation (Commerçon et al. 2011b, 2014, González et al. 2015). In this work, just **grey**
- ✓ See **Lesur et al. (2014)** for a full niMHD implementation in the **PLUTO** code.

$$\begin{aligned}
 \partial_t \rho &+ \nabla \cdot [\rho \mathbf{u}] &= 0 \\
 \partial_t \rho \mathbf{u} &+ \nabla \cdot [\rho \mathbf{u} \otimes \mathbf{u} + P \mathbb{I}] &= -\rho \nabla \Phi - \lambda \nabla E_r + (\nabla \times \mathbf{B}) \times \mathbf{B} \\
 \partial_t E_T &+ \nabla \cdot [\mathbf{u} (E_T + P_T) - \mathbf{B}(\mathbf{B} \cdot \mathbf{u}) + \mathbf{E}_{\text{NIMHD}} \times \mathbf{B}] &= -\rho \mathbf{u} \cdot \nabla \Phi - \mathbb{P}_r \nabla : \mathbf{u} - \lambda \mathbf{u} \nabla E_r + \nabla \cdot \left( \frac{c\lambda}{\rho \kappa_R} \nabla E_r \right) \\
 \partial_t E_r &+ \nabla \cdot [\mathbf{u} E_r] &= -\mathbb{P}_r \nabla : \mathbf{u} + \nabla \cdot \left( \frac{c\lambda}{\rho \kappa_R} \nabla E_r \right) + \kappa_P \rho c (a_R T^4 - E_r) \\
 \partial_t B &- \nabla \times (\mathbf{u} \times \mathbf{B}) - \nabla \times \mathbf{E}_{\text{NIMHD}} &= 0
 \end{aligned}$$

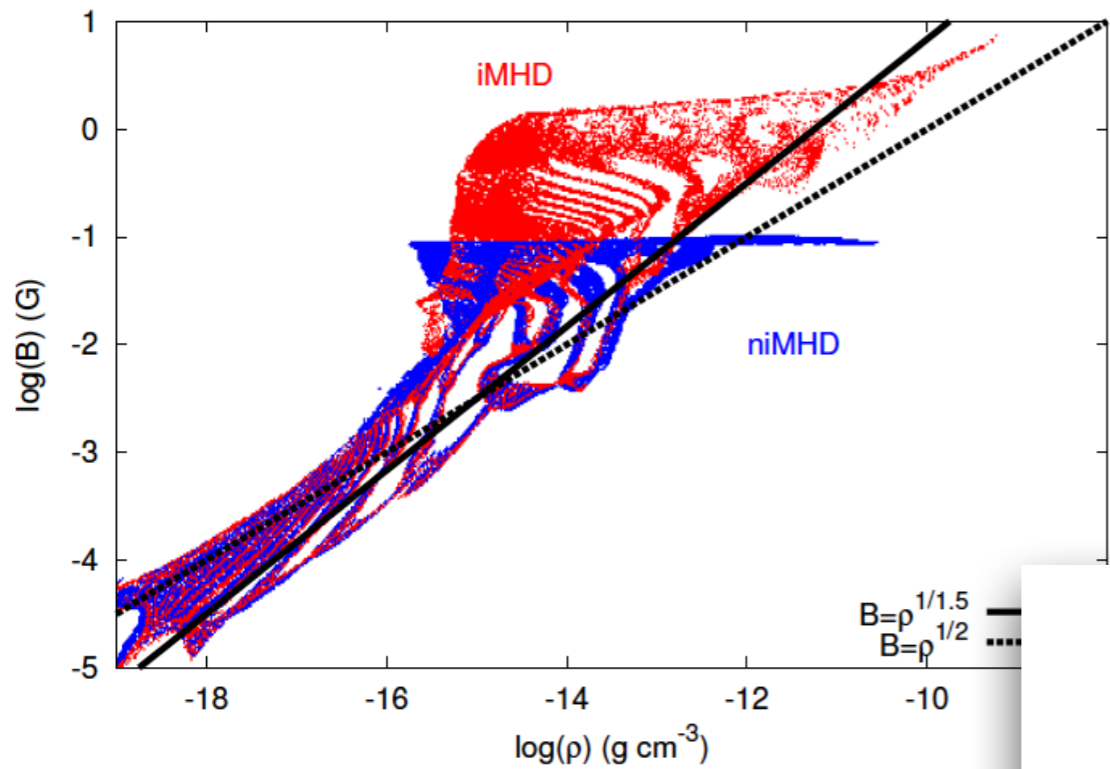
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Gravitational force
Radiative force
Lorentz force

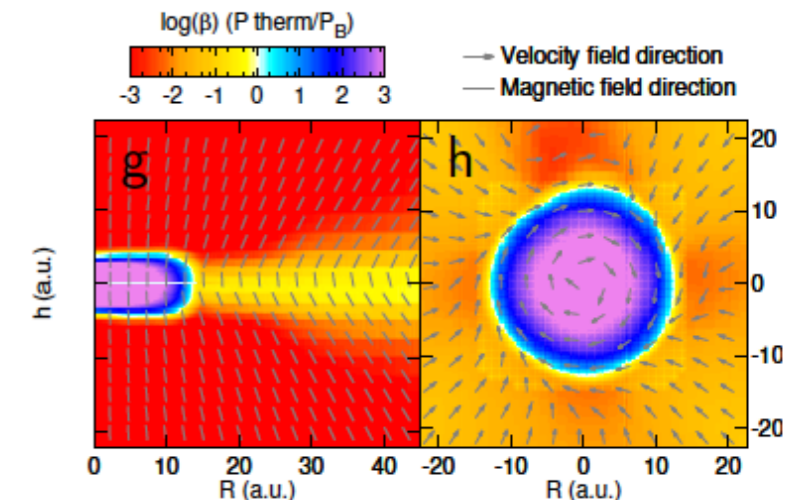
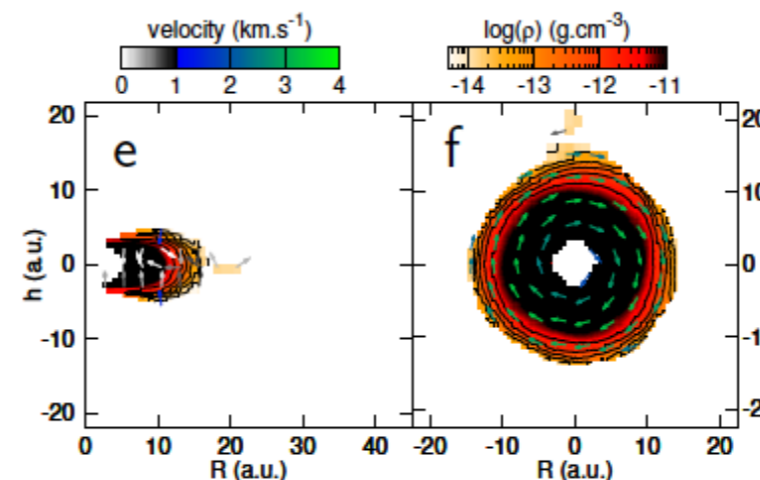
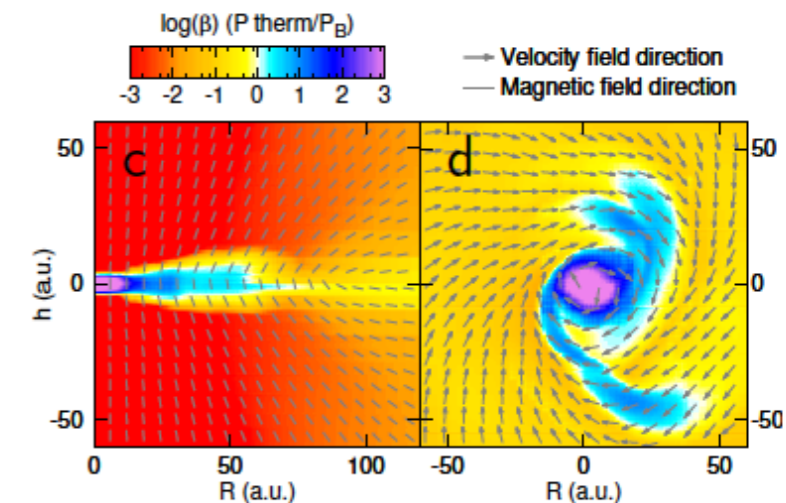
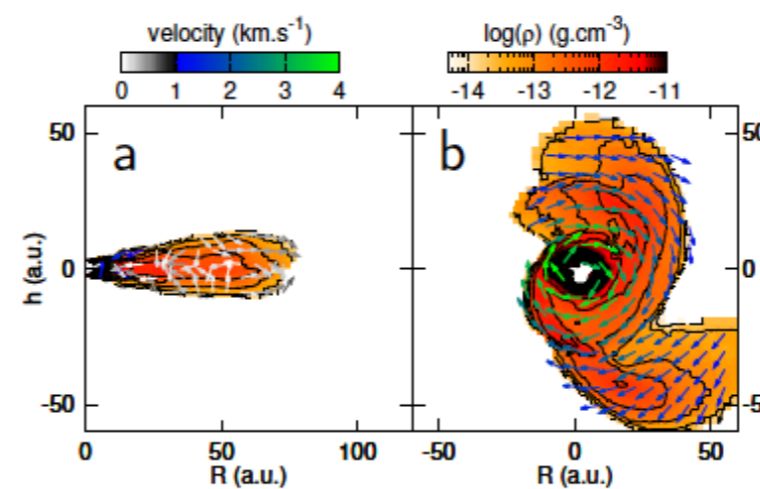
**Resistive EMF**  $\mathbf{E}_{\text{NIMHD}} = +\eta_\Omega \mathbf{J} + \frac{\eta_H}{\|\mathbf{B}\|} \mathbf{J} \times \mathbf{B} - \frac{\eta_{AD}}{\|\mathbf{B}\|^2} \mathbf{J} \times \mathbf{B} \times \mathbf{B}$



# Misalignment & ambipolar diffusion



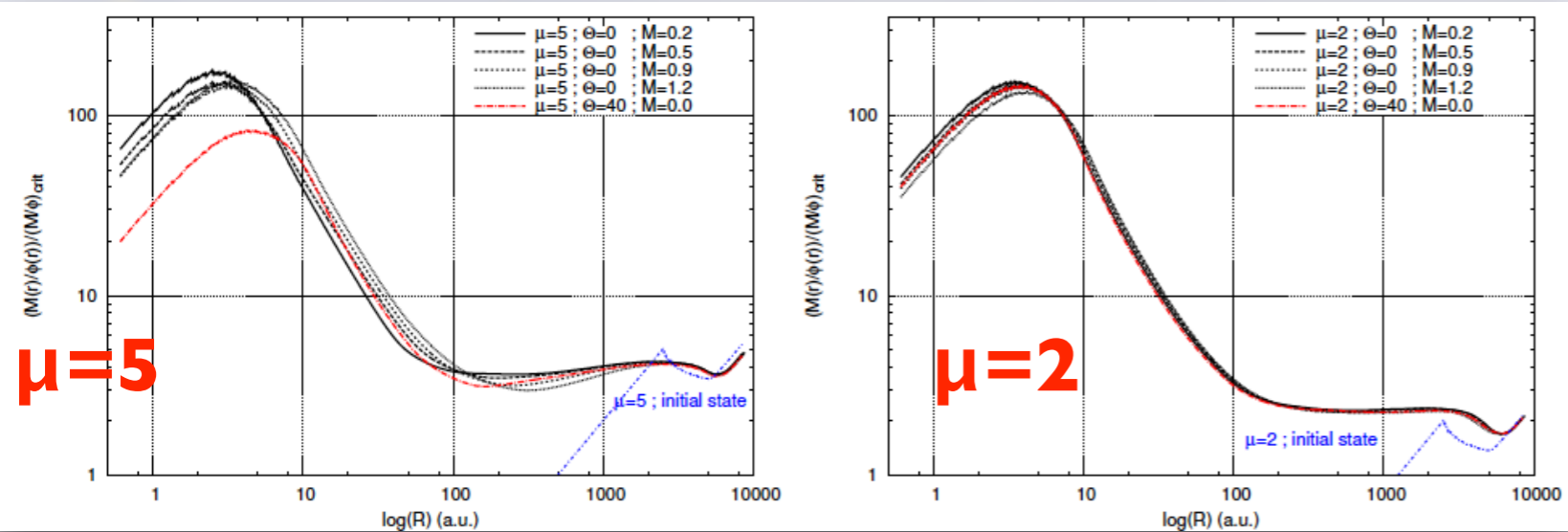
- formation of a **plateau** at  $B \sim 0.1$  G
- **reorganisation** of magnetic field lines (essentially **poloidal**)
- ⇒ **reduced magnetic braking**
- mass and radius of first core do not change
- **weaker outflows** compared to ideal MHD



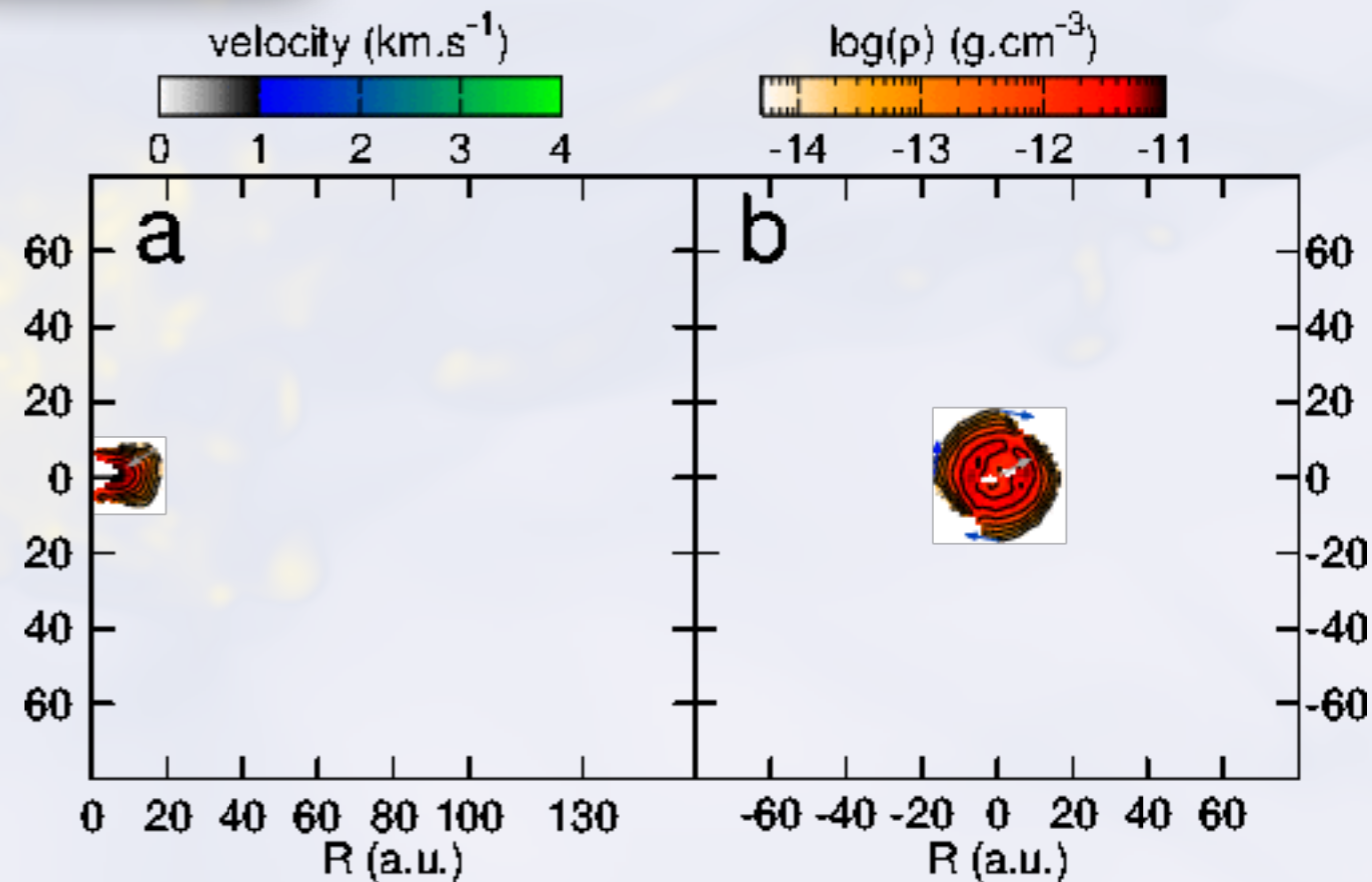
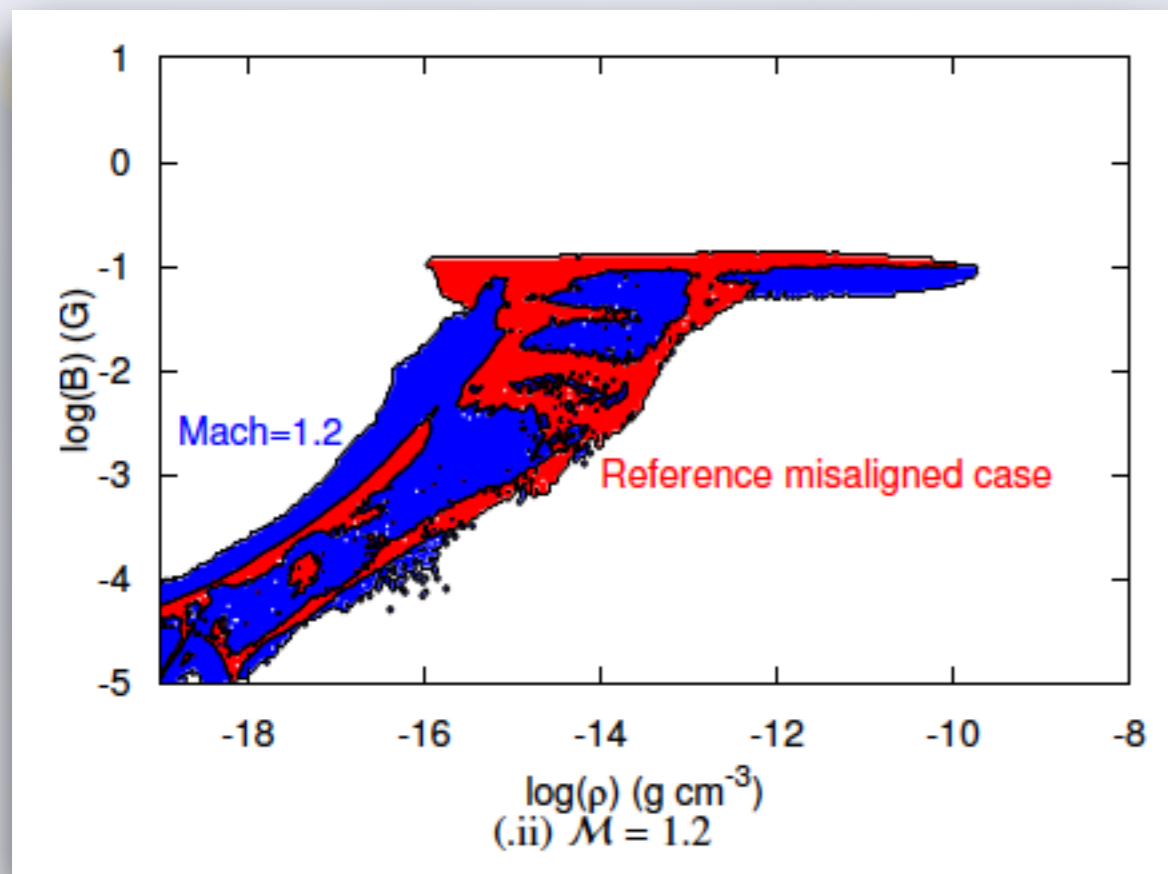
- Rotationally supported disk formation ( $R \sim 50$  AU) - consistent with obs.
- $P_{\text{therm}}/P_{\text{mag}} > 1$  within disks
- **poloidal** magnetic field
- ⇒ initial conditions for protoplanetary disks studies

*Masson et al. 2016*

# Turbulence and ambipolar diffusion



- magnetisation & disc size **does not depend** on turbulence level, nor on the initial magnetic field amplitude



*Commerçon et al. in prep.*

# Magnetically regulated disc size with AD

*Hennebelle et al. (2016)*

$$\tau_{\text{far}} \simeq \frac{B_{\phi} h}{B_z v_{\phi}} \qquad \tau_{\text{br}} \simeq \frac{\rho v_{\phi} 4\pi h}{B_z B_{\phi}}$$

$$\tau_{\text{diff}} \simeq \frac{4\pi h^2}{c^2 \eta_{\text{AD}}} \frac{B_z^2 + B_{\phi}^2}{B_z^2} \simeq \frac{4\pi h^2}{c^2 \eta_{\text{AD}}} \qquad \tau_{\text{rot}} \simeq \frac{2\pi r}{v_{\phi}}$$

$$r_{\text{d,AD}} \simeq 18 \text{ au}$$

$$\times \delta^{2/9} \left( \frac{\eta_{\text{AD}}}{0.1 \text{ s}} \right)^{2/9} \left( \frac{B_z}{0.1 \text{ G}} \right)^{-4/9} \left( \frac{M_{\text{d}} + M_{*}}{0.1 M_{\odot}} \right)^{1/3}$$

- disc size **does not depend** on turbulence level
- weak dependance on the mass

VS.

$$r_{\text{d,hydro}} \simeq \frac{\Omega_0^2 R_0^4}{4\pi/3\rho_0 R_0^3 G} = 3\beta R_0 = 106 \text{ AU} \frac{\beta}{0.02} \left( \frac{M}{0.1 M_{\odot}} \right)^{1/3} \left( \frac{\rho_0}{10^{-18} \text{ g cm}^{-3}} \right)^{-1/3}$$

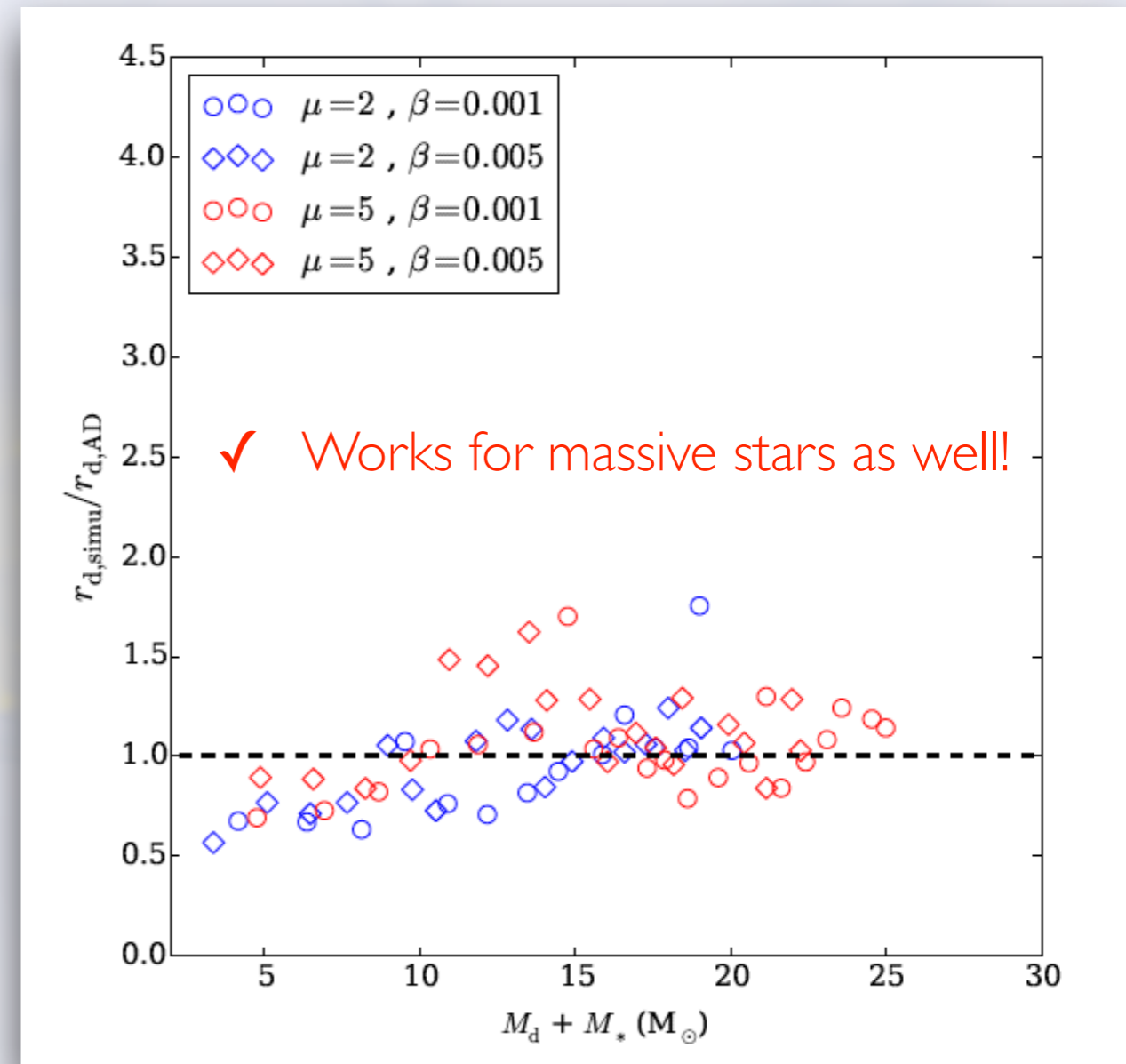
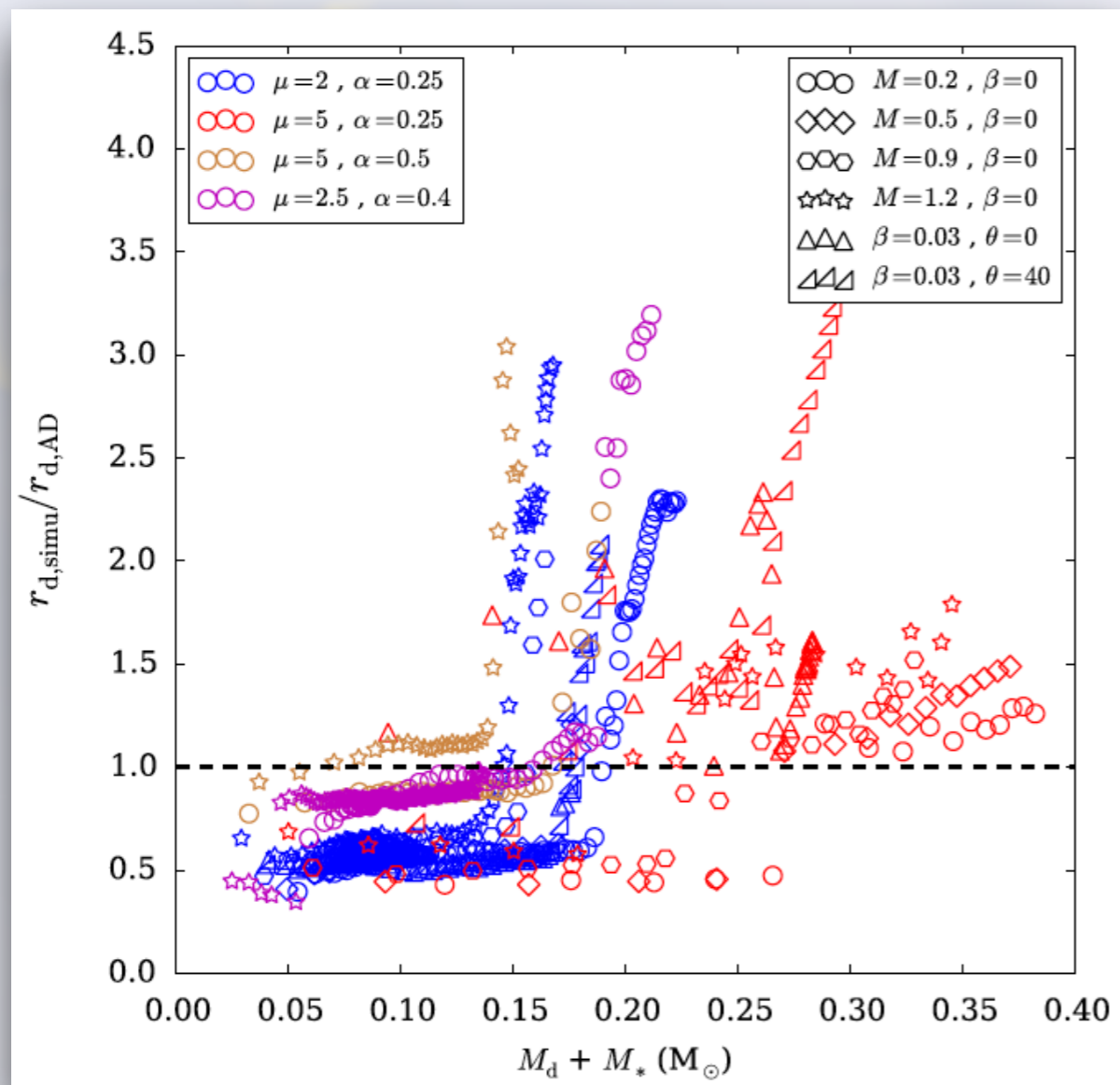


# Magnetically regulated disc size with AD

$$r_{d,AD} \simeq 18 \text{ au}$$

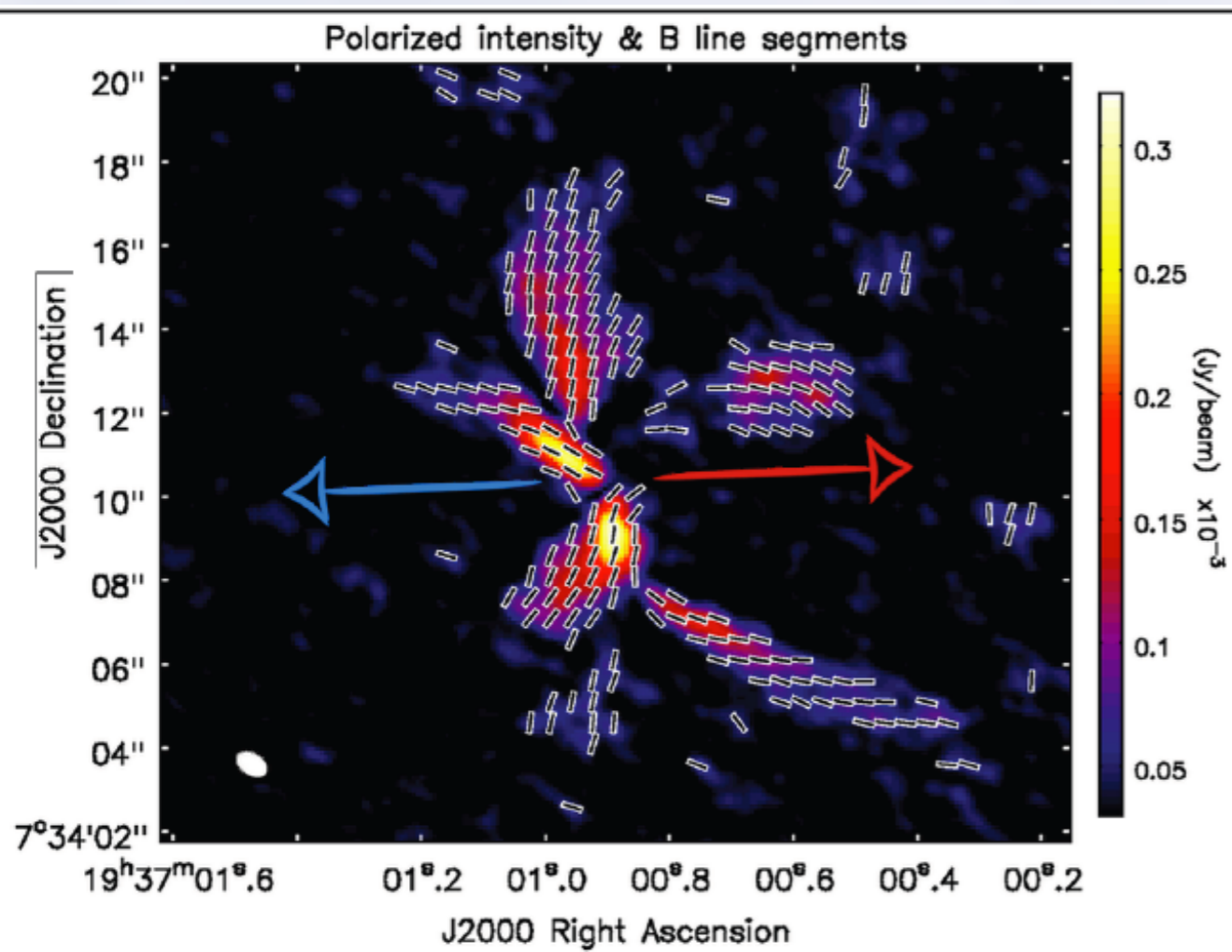
$$\times \delta^{2/9} \left( \frac{\eta_{AD}}{0.1 \text{ s}} \right)^{2/9} \left( \frac{B_z}{0.1 \text{ G}} \right)^{-4/9} \left( \frac{M_d + M_*}{0.1 M_\odot} \right)^{1/3}$$

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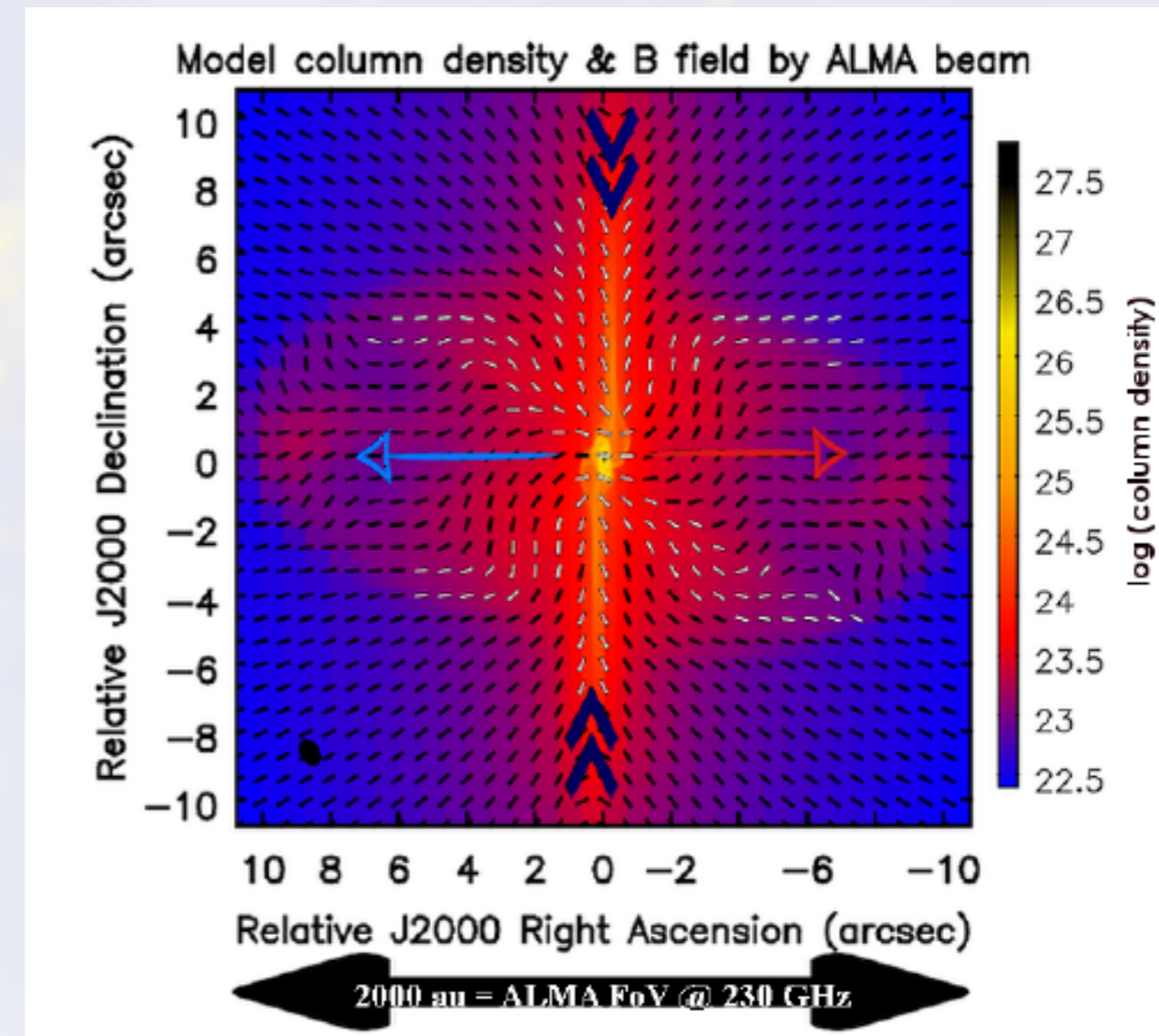
# Magnetically regulated collapse?



Yen et al. (2015): No Keplerian disk with  $R > 10$  AU  
=> Magnetic braking or young age?

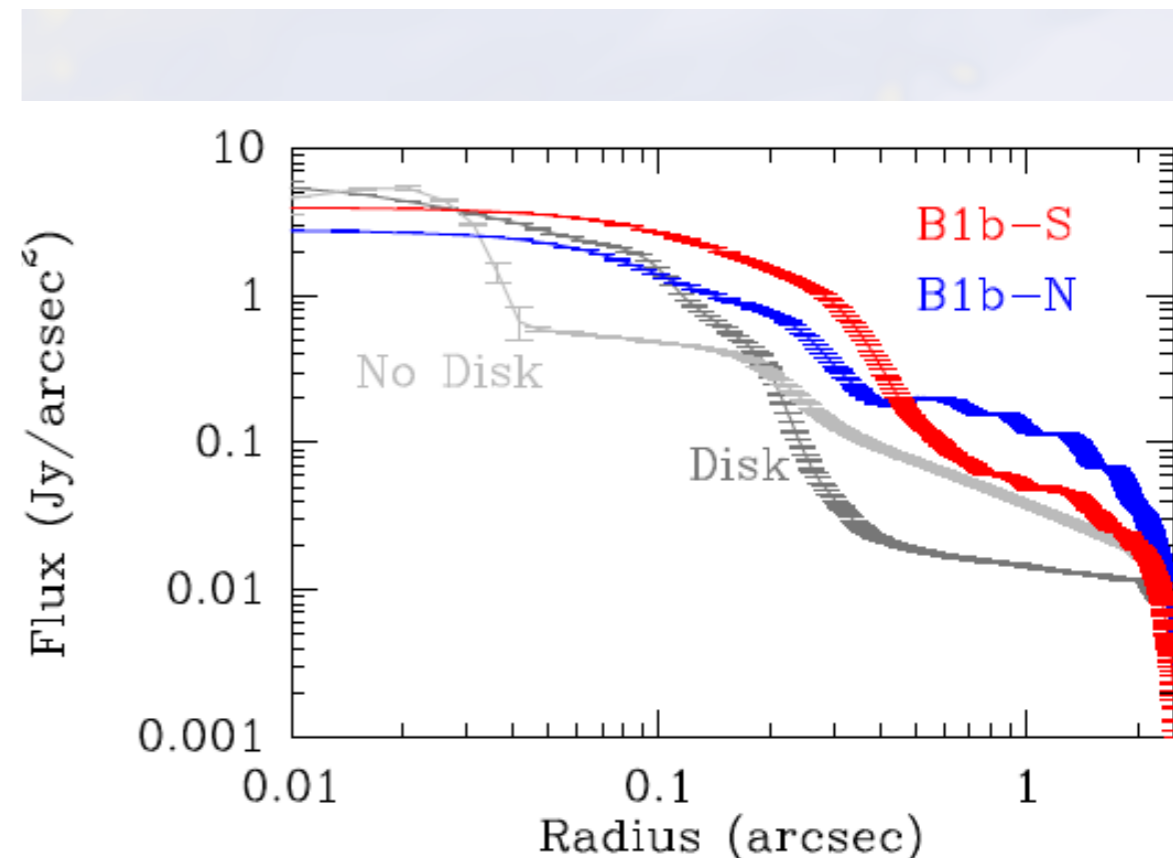
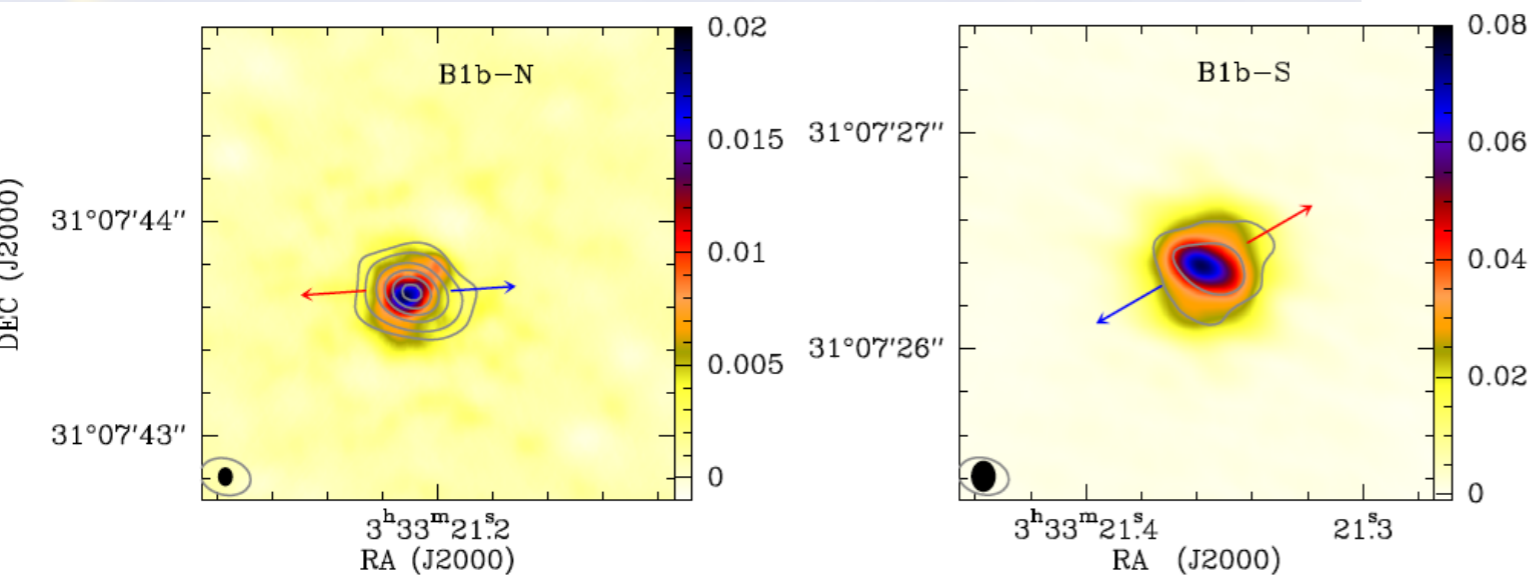
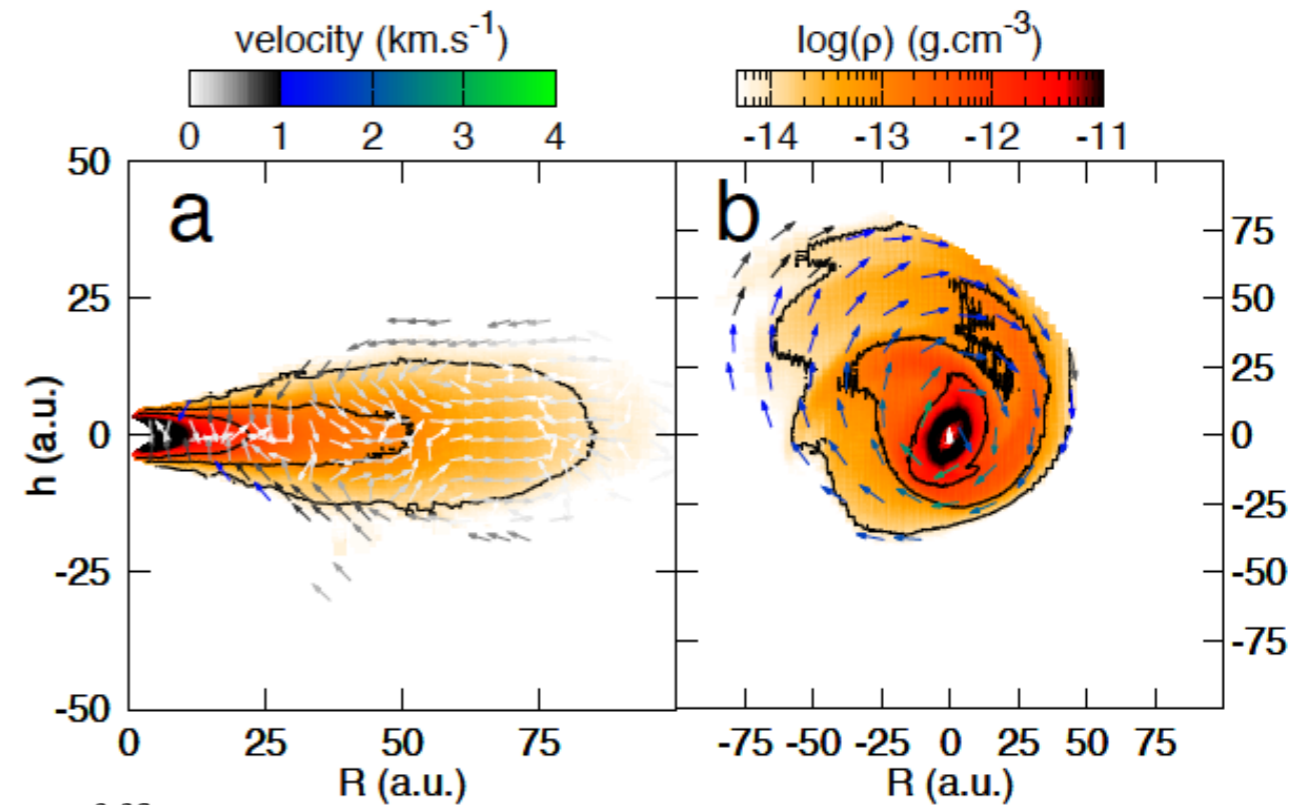
## Maury et al. (2018)

- Source B335 Class 0
- ALMA 0.8" @ 233 GHz ( $\sim 60$  AU)
- data "consistent" with magnetised collapse model ( $2.5M_{\odot}$ ;  $\mu=6$ ; low rotation)



# Comparison with ALMA observations

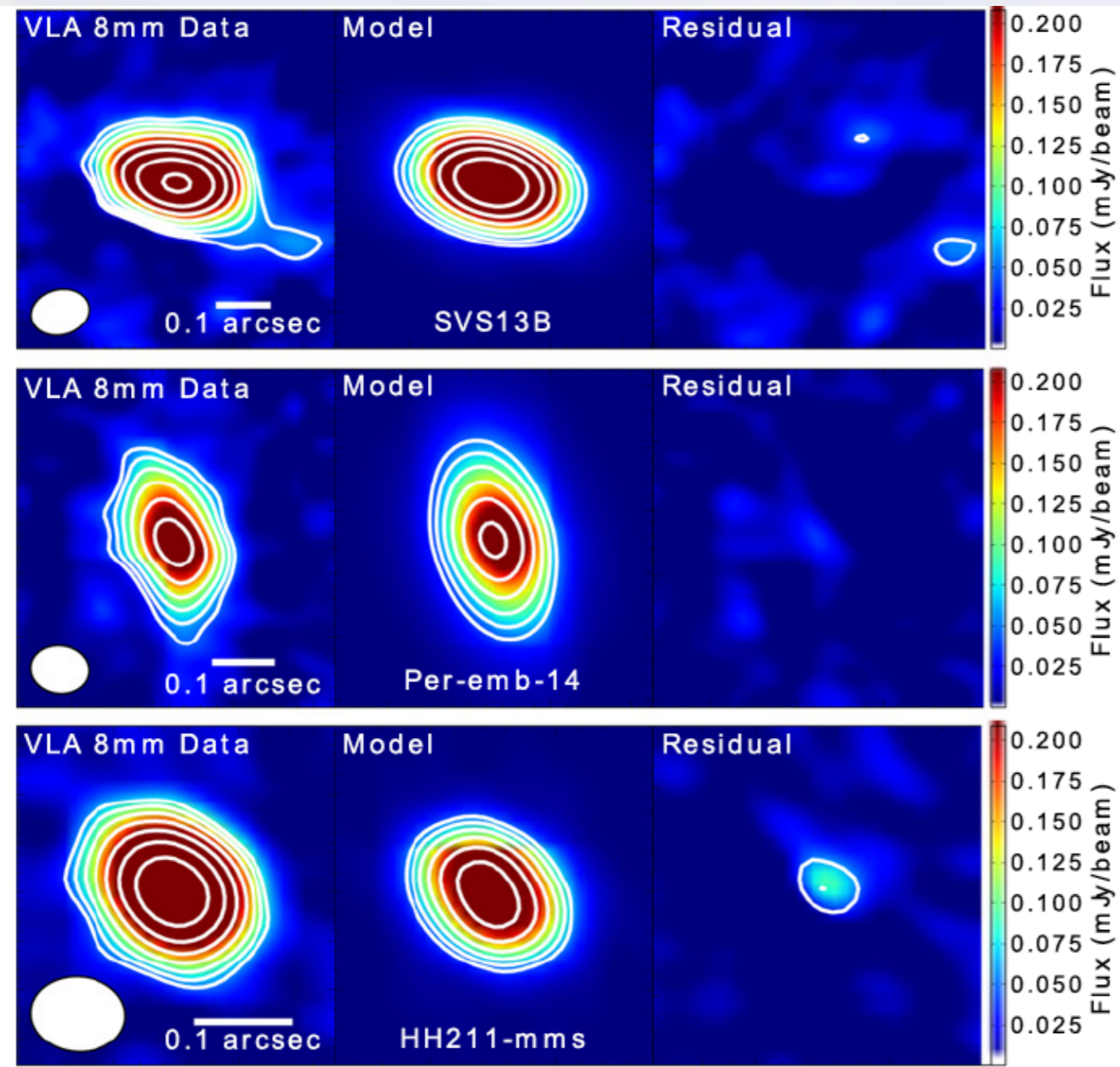
- Source Barnard 1b-N: first core candidate
- ALMA 0.06" @ 350 GHz (~15 AU)
- data "consistent" with collapse model ( $1M_{\odot}$ ;  $\mathcal{M}=1.2$ ;  $\mu=2$ )
- data compatible with disc growth with time



*Gerin et al. A&A (2017)*



# Evidence of disks at class 0 & I stages



**Table 3**  
8mm Best-fit Modeling Results

Source	$q$	$\gamma$	$R_c$ (AU)	$\chi^2_{\text{reduced}}$
SVS13B	0.25	$0.21^{+0.23}_{-0.20}$	$24.28^{+2.1}_{-1.7}$	2.194
	0.50	$0.42^{+0.25}_{-0.21}$	$25.50^{+1.9}_{-1.5}$	2.185
	0.75	$0.63^{+0.24}_{-0.22}$	$26.46^{+1.6}_{-1.4}$	2.175
	1.00	$0.85^{+0.26}_{-0.23}$	$27.28^{+1.4}_{-1.2}$	2.164
class I	0.25	$0.08^{+0.02}_{-0.16}$	$21.9^{+0.8}_{-0.9}$	1.556
	0.50	$0.26^{+0.15}_{-0.17}$	$23.3^{+1.1}_{-1.0}$	1.558
	0.75	$0.44^{+0.16}_{-0.17}$	$24.6^{+1.4}_{-1.1}$	1.560
	1.00	$0.64^{+0.16}_{-0.18}$	$25.7^{+1.4}_{-1.3}$	1.563
Per-emb-14	0.25	$-0.11^{+0.16}_{-0.00}$	$28.5^{+2.3}_{-2.1}$	1.110
	0.50	$0.09^{+0.08}_{-0.21}$	$30.6^{+2.8}_{-2.3}$	1.114
	0.75	$0.27^{+0.17}_{-0.24}$	$32.5^{+2.2}_{-2.8}$	1.119
	1.00	$0.48^{+0.19}_{-0.23}$	$33.9^{+3.6}_{-3.1}$	1.123
Per-emb-30	0.25	$0.02^{+0.18}_{-0.31}$	$14.0^{+1.0}_{-0.9}$	1.100
	0.50	$0.20^{+0.04}_{-0.32}$	$14.9^{+1.9}_{-1.1}$	1.102
	0.75	$0.39^{+0.30}_{-0.34}$	$15.8^{+1.9}_{-1.3}$	1.104
	1.00	$0.59^{+0.14}_{-0.33}$	$16.5^{+31.3}_{-1.6}$	1.107
HH211-mms	0.25	$0.48^{+0.40}_{-0.78}$	$10.5^{+0.8}_{-0.8}$	1.009
	0.50	$0.65^{+0.43}_{-0.82}$	$11.0^{+1.0}_{-0.9}$	1.009
	0.75	$0.81^{+0.42}_{-0.79}$	$11.5^{+1.2}_{-1.2}$	1.009
	1.00	$1.01^{+0.44}_{-0.81}$	$11.9^{+1.4}_{-1.3}$	1.009
IC348 MMS	0.25	$-0.58^{+0.11}_{-0.11}$	$25.7^{+2.8}_{-2.2}$	1.085
	0.50	$-0.39^{+0.19}_{-0.11}$	$29.0^{+3.2}_{-2.6}$	1.096
	0.75	$-0.19^{+0.11}_{-0.27}$	$31.6^{+4.1}_{-2.9}$	1.107
	1.00	$0.02^{+0.07}_{-0.11}$	$33.7^{+4.3}_{-3.1}$	1.118
Per-emb-8	0.25	$0.01^{+0.16}_{-0.19}$	$19.0^{+1.2}_{-1.1}$	1.099
	0.50	$0.20^{+0.17}_{-0.20}$	$20.2^{+1.4}_{-1.3}$	1.107
	0.75	$0.40^{+0.17}_{-0.21}$	$21.2^{+1.6}_{-1.4}$	1.114
	1.00	$0.61^{+0.17}_{-0.20}$	$22.1^{+1.8}_{-1.6}$	1.122

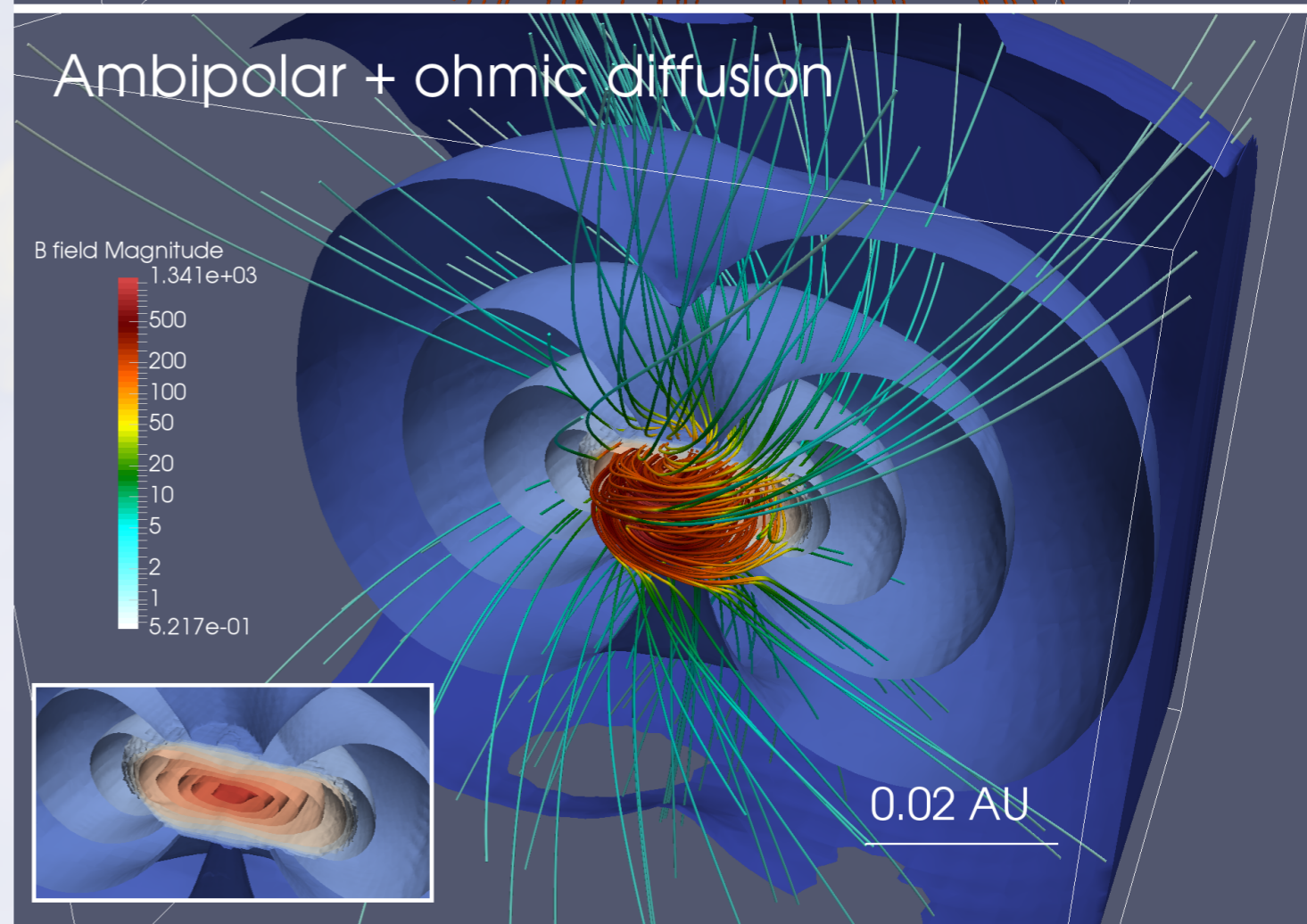
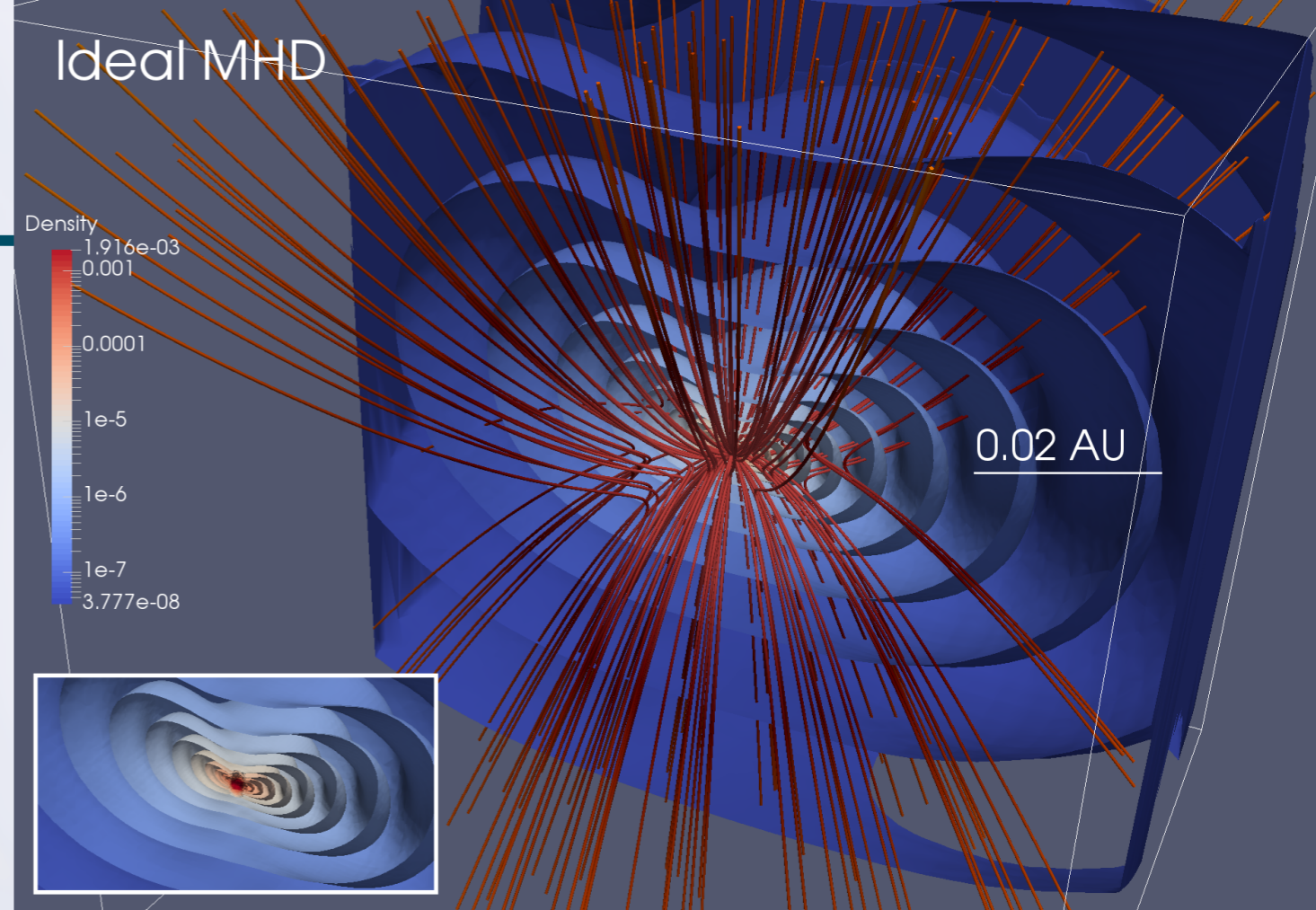
*Segura-Cox et al. (2016)*



# Second collapse

- Ohmic + ambipolar diffusion
  - non-ideal gas EOS
- Saumon, Chabrier & Von Horn (1995)*
- maximum resolution :  $\Delta x \sim 8 \times 10^{-5}$  AU (21 AMR levels)
  - Comparaison ideal MHD (80 000 hCPU) vs. non ideal MHD (180 000 hCPU)

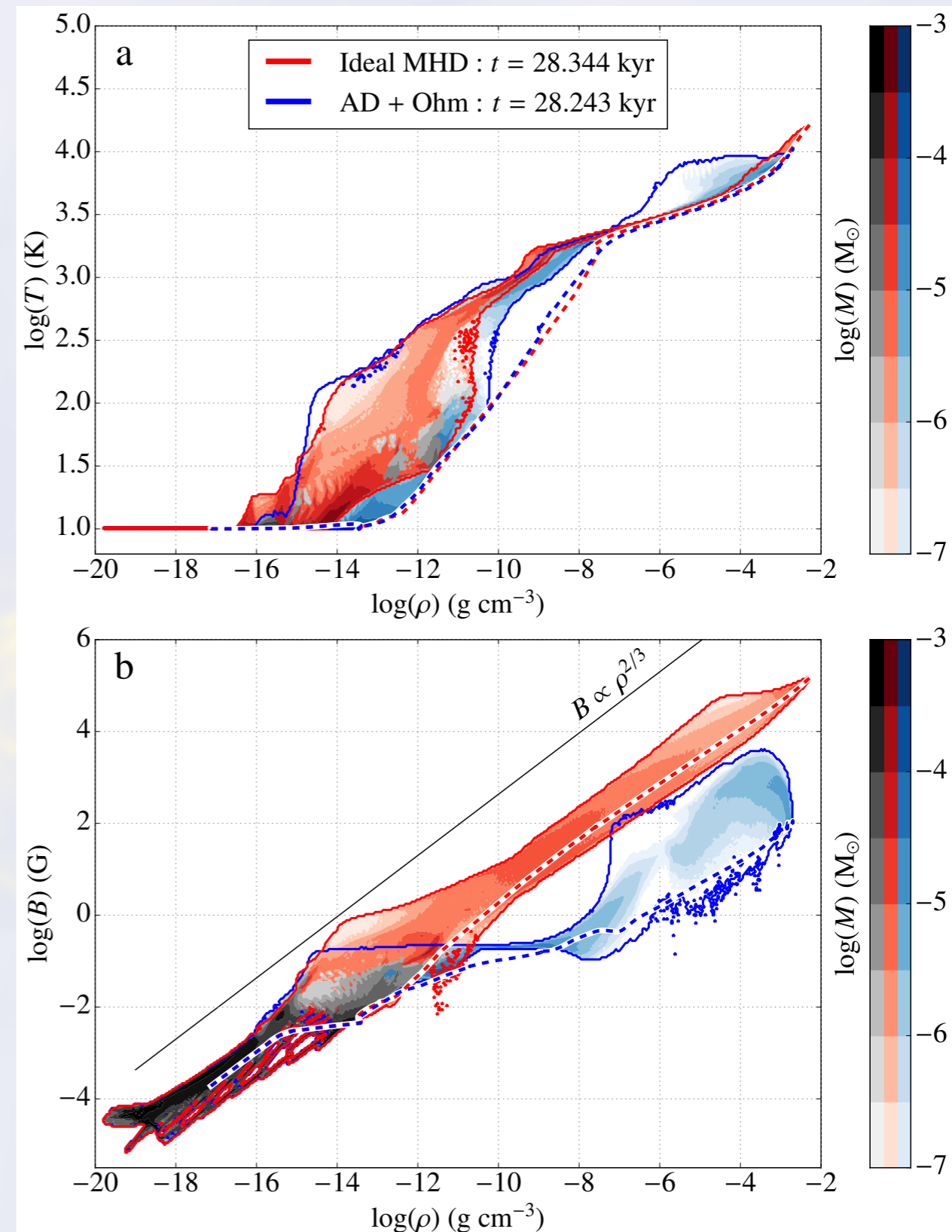
*Vaytet et al. (2018)*





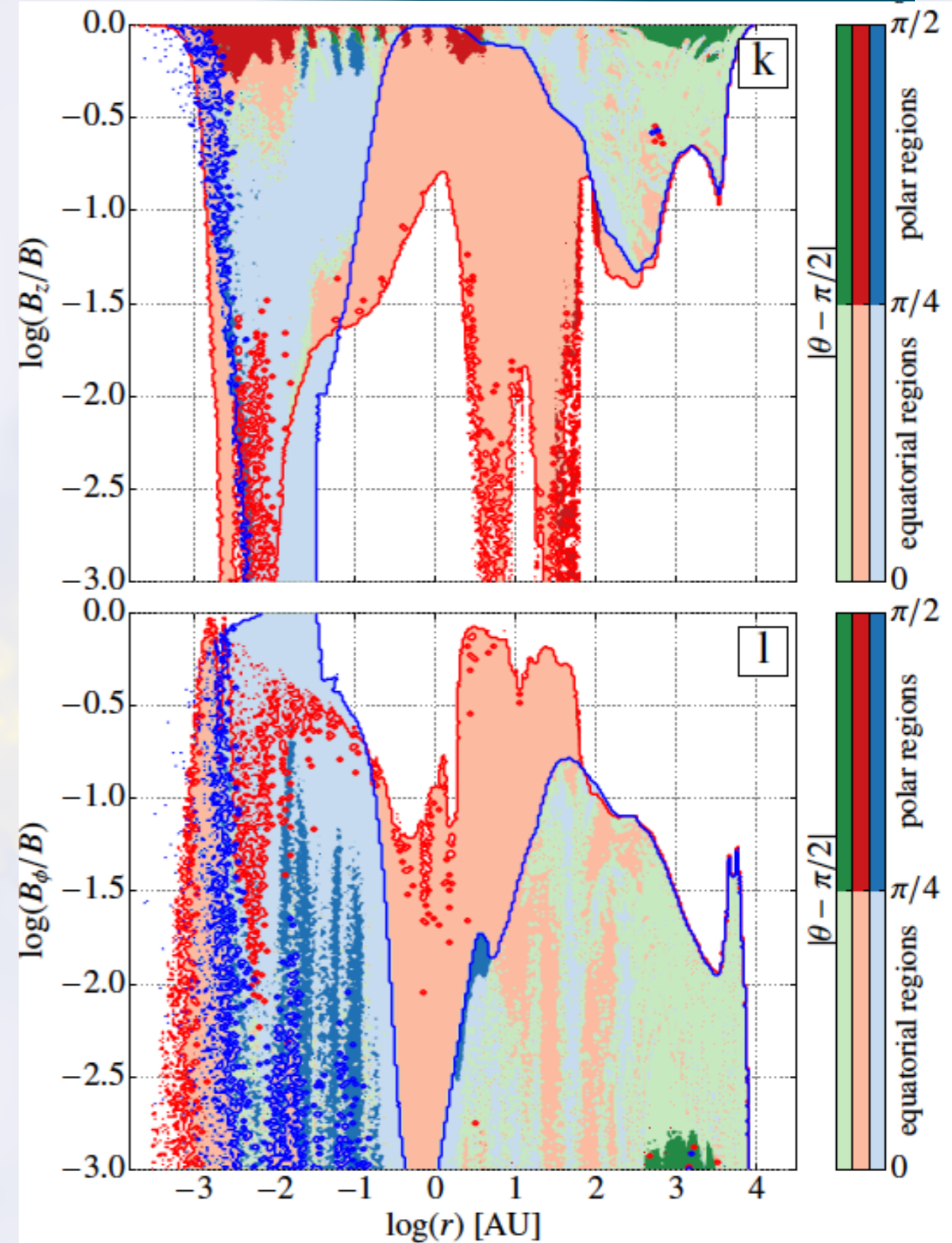
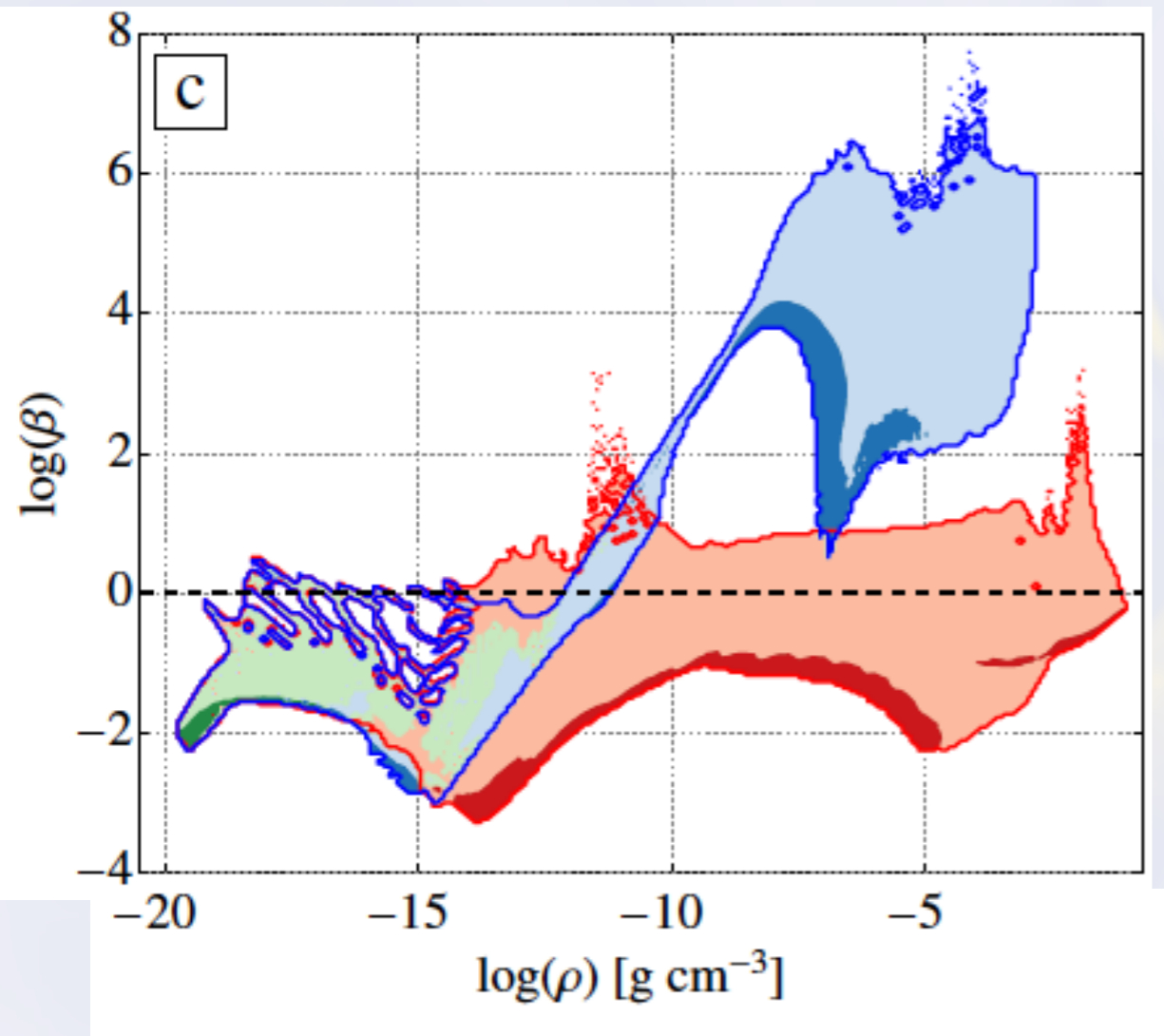
# Second collapse

Magnetic flux reduced by  $\sim 3$  orders of magnitude only with **ambipolar diffusion** and **Ohmic diffusion**



# Second collapse

- ideal MHD
- non-ideal MHD



*Vaytet et al. (2018)*





# Properties of the radiative accretion shock

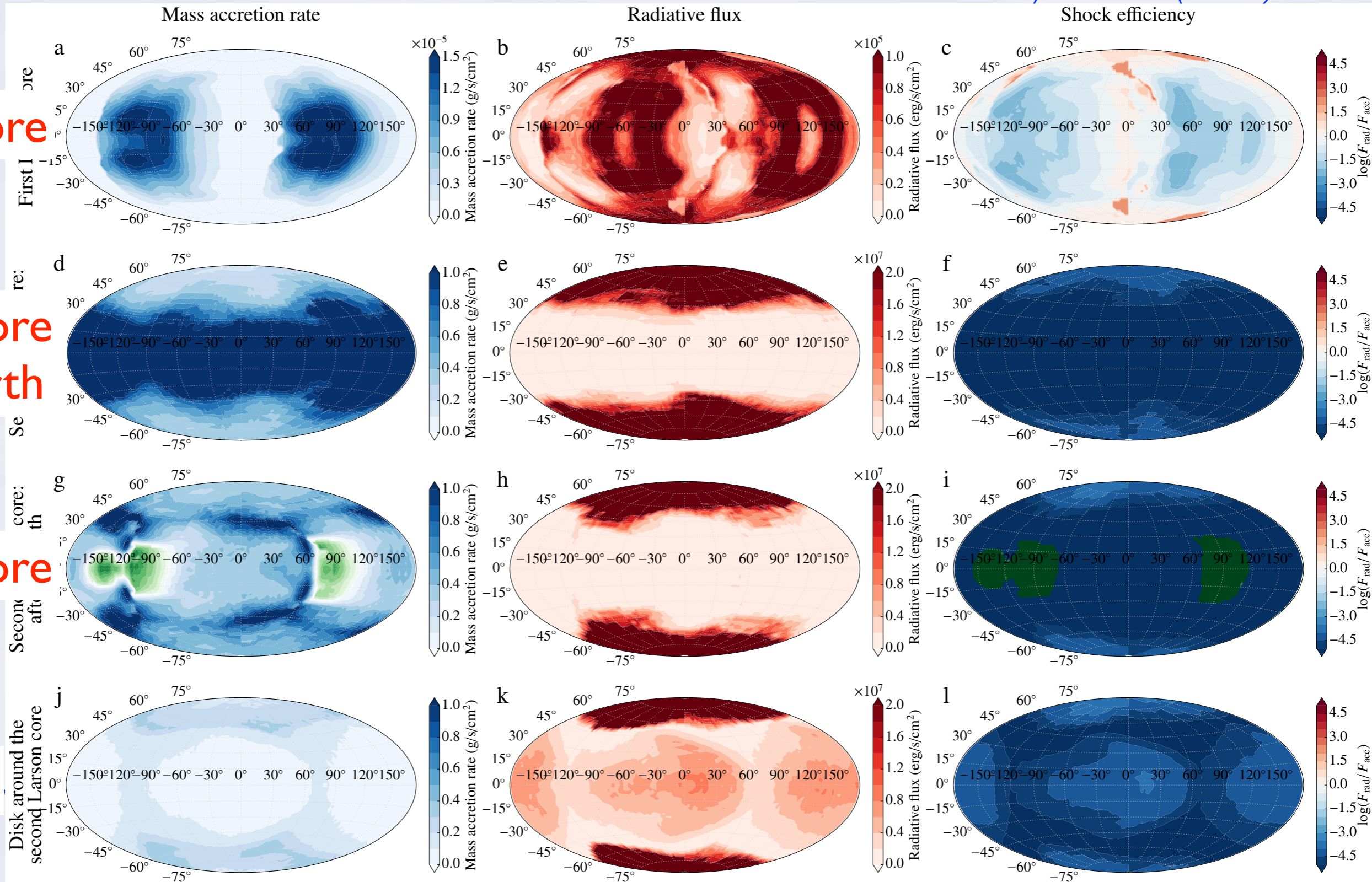
Vaytet et al. (2018)

1<sup>st</sup> core

2<sup>nd</sup> core  
@birth

2<sup>nd</sup> core

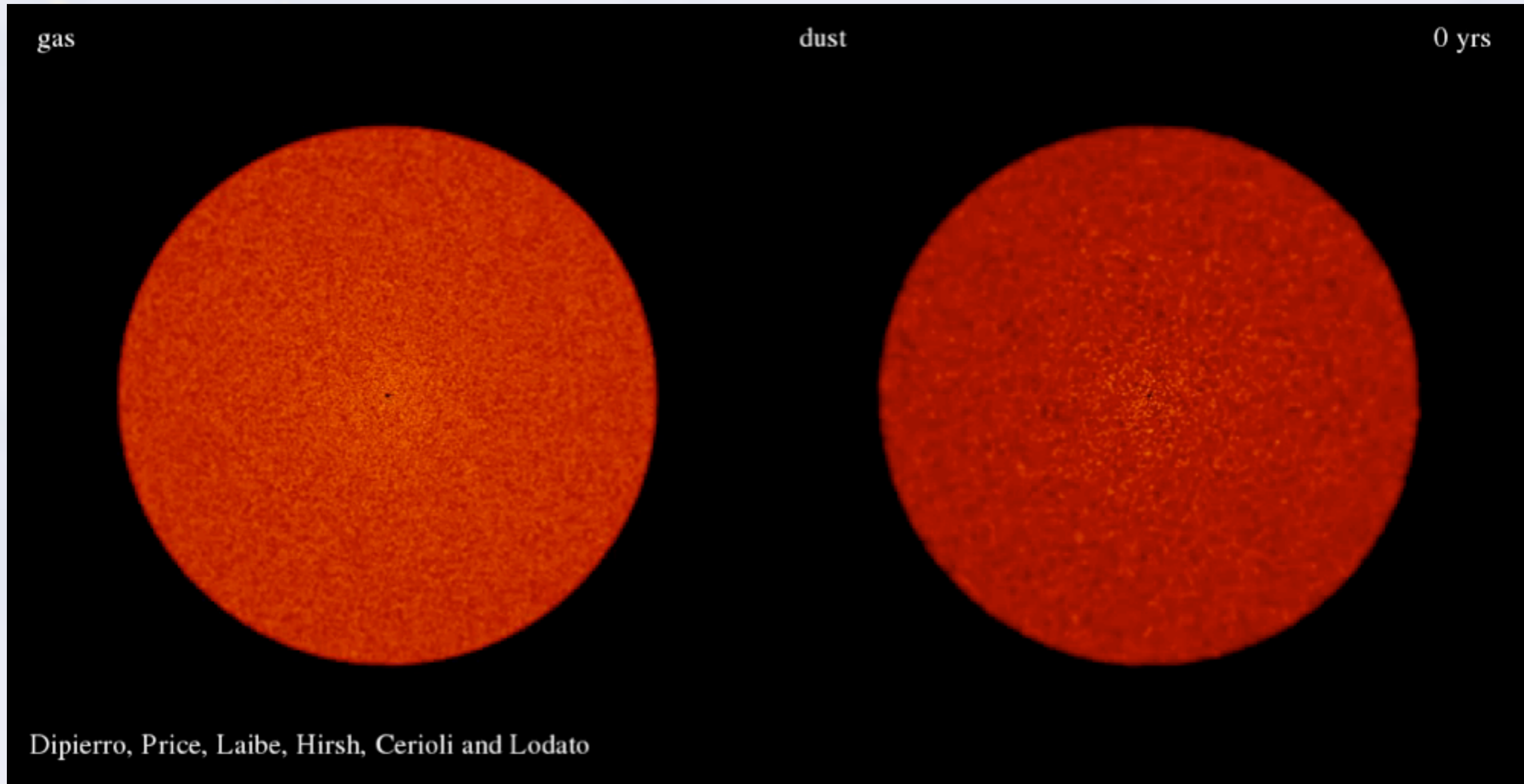
Disc





# What is next?

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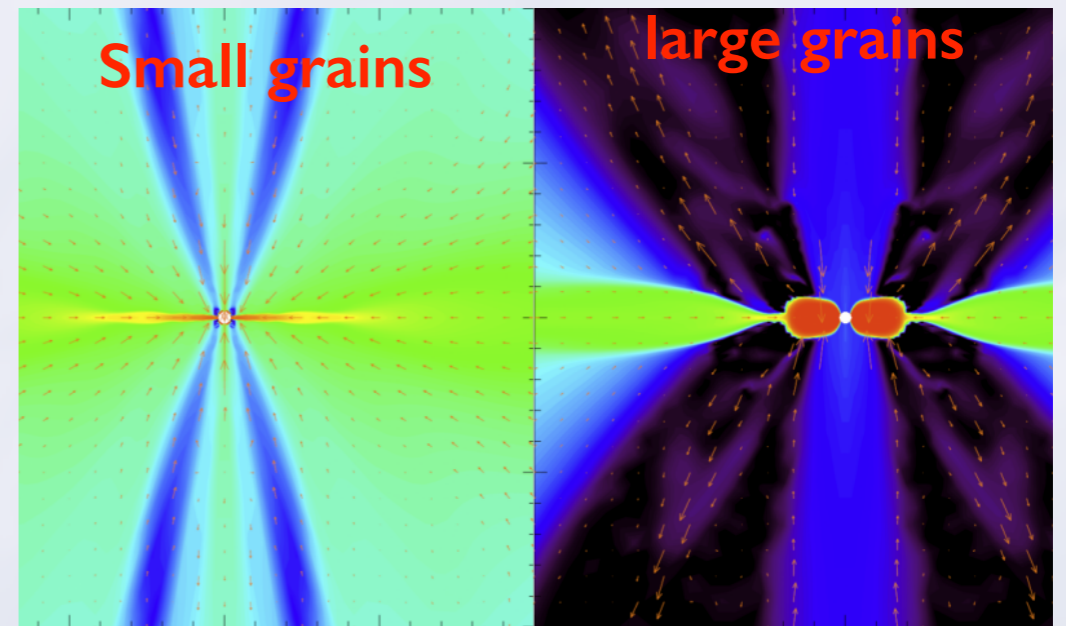
*Dipierro et al. (2015)*

# What is next?

- **Follow the dust dynamics at all scales**

- dust growth
- dust charge
- gas-to-dust ratio

*Zhao et al (2016)*



- **Couple the processes**

- magneto hydrodynamics: chemistry + dust + magnetic resistivities
- radiation hydrodynamics: chemistry + dust + opacities
- track cosmic rays ionisation

- **Couple the scales**

- galaxy evolution to molecular clouds (e.g., Renaud et al. 2013)
- self-regulated ISM, from diffuse ISM to collapsing dense cores (Hennebelle et al.)
- protoplanetary disc evolution with accreting envelop

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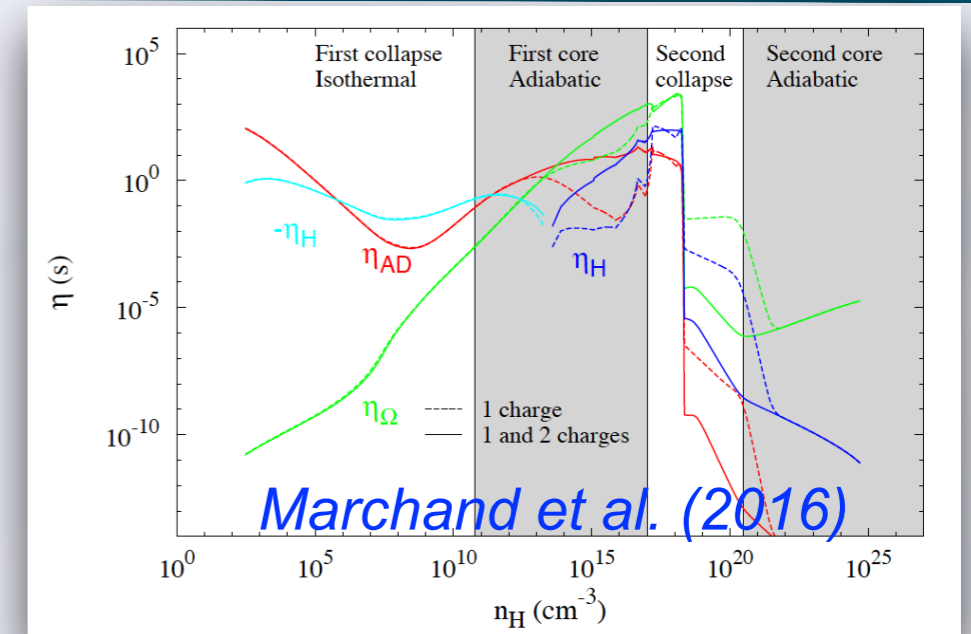
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# Conclusion

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- ☑ Discs are back! Disc size regulated by ambipolar diffusion
  - ☑ Magnetic fields **cannot be neglected. Non-ideal MHD** as well...
  - ☑ Magnetised models compare well with observations
- 
- ➡ Dust evolution
  - ➡ Second collapse - parameter study
  - ➡ Massive stars



# Astrosim Conf, 8-11 octobre 2018, Lyon

ASTRO  
SIM

## Highlights and prospects for numerical astrophysics in France

8-11 Oct 2018 Lyon (France)

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### MAIN MENU

Home

Registration

Program

List of Participants

### HELP

@ Contact

### Rational

High Performance Computing (HPC) has become an essential tool of all fields of astronomy and astrophysics. In particular, numerical simulations have become the tool of choice to carry out quantitative theoretical investigations, which take into account the full complexity of astrophysical objects and their contexts, and reach a level of precision which matches the ever more demanding level of detail reached with observations. Today, numerical simulations of astrophysical objects have become the principal vector of the dialog between observations and theory, and they have reached a quality which allows us to use them for a direct interpretation of observations (from large-scale cosmological surveys to observations of atmospheric phenomena on the surface of giant planets).

As observations, simulations require high-performance instruments: super-computers. In less than 10 years, there has been a spectacular increase of the gigantic computing power that is now available to researchers, both at the national (through GENCI) and european level (through PRACE). This is accompanied by profound evolutions of the methods used by specialists of numerical simulations. More and more research projects are now performed in close collaborations with software engineers. With exascale computers coming soon, such a professionalism of high performance computing (HPC) is expected to become more and more frequent. Like the other communities in France, astrophysicists must prepare to that transition.

The aim of this workshop is to bring together researchers working in France in all fields of numerical astrophysics, along with representatives of the HPC infrastructure at the regional, national, and European level, and key players in the forthcoming ExaScale HPC. Invited presentations from astronomers will both highlight astrophysical results and propose a reflexion on the evolutions of their work. Other invited presentations will present the expected evolution of the HPC landscape in France and Europe, both in terms of organisation and of technologies. A large number of contributed talks will complement the program.

### Confirmed Speakers

Dominique Aubert (ObAS), Edouard Audit (MdS), Frédéric Bournaud (CEA), Allan Sacha Brun (CEA), Fabien Casse (APC), Benoît Cerutti (IPAG), Boris Dintrons (CINES), Yohan Dubois (IAP), François Forget (LMD), Philippe Grandclément (LUTH), Oliver Hahn (OCA), Patric Hennebelle (CEA), Laurène Jouve (IRAP), Guillaume Laibe (CRAL), Geoffroy Lesur (IPAG), Sophie Masson (LESIA), Stéphane Mazevet (LUTH), Patrick Michel (OCA), Michel Rieutord (IRAP), Roch Smets (UPMC), Pascal Tremblin (CEA)